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**ON SOLUTION OF THE INVERSE
STURM-LIOUVILLE PROBLEM WITH
DISCONTINUOUS COEFFICIENTS**

Abstract

In the paper we prove a theorem on necessary and sufficient conditions on solvability of the inverse problem of spectral analysis for Sturm-Liouville operator with discontinuous coefficients and give solution algorithm of the inverse problem.

Consider the boundary value problem

$$-y'' + q(x)y = \lambda^2 \rho(x)y, \quad 0 \leq x \leq \pi, \tag{1}$$

$$y(0) = 0, \quad y(\pi) = 0, \tag{2}$$

where $q(x) \in L_2(0, \pi)$ is a real-valued function, λ is a complex parameter and $\rho(x)$ is a piecewise-constant function:

$$\rho(x) = \begin{cases} 1, & 0 \leq x \leq a, \\ \alpha^2, & a < x \leq \pi, \end{cases} \quad 0 < \alpha \neq 1. \tag{3}$$

We assume that $a(1 + \alpha) > \pi\alpha$.

The following theorem gives necessary and sufficient conditions for solvability of the inverse problem of spectral analysis for boundary value problem (1) – (3).

Theorem. *For the sequences $\{\lambda_n^2, \alpha_n\}_{n \geq 1}$, where $\lambda_n \neq \lambda_m$ for $n \neq m$, $\alpha_n > 0$ for all n to be the spectral data of a problem $L(q(x))$ of the form (1) – (3) with $q(x) \in L_2(0, \pi)$, it is necessary and sufficient to satisfy conditions*

$$\lambda_n = \lambda_n^0 + \frac{d_n}{\lambda_n^0} + \frac{k_n}{n}, \quad \alpha_n = \alpha_n^0 + \frac{t_n}{n}. \tag{4}$$

Here λ_n^0 are the zeros of the function

$$\Delta_0(\lambda) \equiv \frac{1}{2} \left(1 + \frac{1}{\alpha} \right) \frac{\sin \lambda (\alpha\pi - \alpha a + a)}{\lambda} + \frac{1}{2} \left(1 - \frac{1}{\alpha} \right) \frac{\sin \lambda (-\alpha\pi + \alpha a + a)}{\lambda};$$

$$\alpha_n^0 = \int_0^\pi s_0^2(x, \lambda_n^0) \rho(x) dx;$$

$$s_0(x, \lambda) = \frac{1}{2} \left(1 + \frac{1}{\sqrt{\rho(x)}} \right) \frac{\sin \lambda \mu^+(x)}{\lambda} + \frac{1}{2} \left(1 - \frac{1}{\sqrt{\rho(x)}} \right) \frac{\sin \lambda (\mu^-(x))}{\lambda};$$

$$\mu^\pm(x) = \pm x \sqrt{\rho(x)} + a \left(1 \mp \sqrt{\rho(x)} \right);$$

d_n is a bounded sequence; $\{k_n\} \in l_2$, $\{t_n\} \in l_2$.

Proof. Necessity of the theorem is proved in [1]. Let's prove sufficiency. Let the real numbers $\{\lambda_n^2, \alpha_n\}_{n \geq 1}$ of the form (4) be given. Construct the functions $F_0(x, t)$ and $F(x, t)$ by the formulas

$$F_0(x, t) = \sum_{n=1}^{\infty} \left(\frac{s_0(t, \lambda_n) \sin \lambda_n x}{\alpha_n \lambda_n} - \frac{s_0(t, \lambda_n^0) \sin \lambda_n^0 x}{\alpha_n^0 \lambda_n^0} \right) \quad (5)$$

and

$$F(x, t) = \frac{1}{2} \left(1 + \frac{1}{\sqrt{\rho(x)}} \right) F_0(\mu^+(x), t) + \frac{1}{2} \left(1 - \frac{1}{\sqrt{\rho(x)}} \right) F_0(\mu^-(x), t). \quad (6)$$

Let $A(x, t)$ be a solution of the main equation of the inverse problem for boundary value problem (1) – (3) (see [2]):

$$\begin{aligned} & \frac{2}{1 + \sqrt{\rho(t)}} A(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A(x, 2a-t) + F(x, t) + \\ & + \int_0^{\mu^+(x)} A(x, \xi) F_0(\xi, t) d\xi = 0, \quad 0 < t < x. \end{aligned} \quad (7)$$

Construct the function $s(x, \lambda)$ by the formula

$$s(x, \lambda) = s_0(x, \lambda) + \int_0^{\mu^+(x)} A(x, t) \frac{\sin \lambda t}{\lambda} dt, \quad (8)$$

and the function $q(x)$ by the formula

$$q(x) = \frac{4\rho(x)}{\sqrt{\rho(x)} + 1} \frac{d}{dx} A(x, \mu^+(x)). \quad (9)$$

Denote

$$b(x) = \sum_{n=1}^{\infty} \left(\frac{\cos \lambda_n x}{\alpha_n \lambda_n^2} - \frac{\cos \lambda_n^0 x}{\alpha_n \lambda_n^0} \right).$$

We can show that $b(x) \in W_2^1(0, \pi)$,

$$\begin{aligned} F_0(x, t) &= \frac{1}{4} \left(1 + \frac{1}{\sqrt{\rho(t)}} \right) [b(x - \mu^+(t)) - b(x + \mu^+(t))] + \\ &+ \frac{1}{4} \left(1 - \frac{1}{\sqrt{\rho(t)}} \right) [b(x - \mu^-(t)) - b(x + \mu^-(t))], \end{aligned}$$

and the functions $F_0(x, t)$, $F(x, t)$, $A(x, t)$ have the same smoothness.

According to (5) and (6) we have

$$F_{0tt}''(x, t) = \rho(t) F_{0xx}''(x, t), \quad \rho(t) F_{xx}''(x, t) = \rho(x) F_{tt}''(x, t), \quad (10)$$

$$F_0(x, t)|_{x=0} = 0, \quad F_0(x, t)|_{t=0} = 0, \quad (11)$$

$$\frac{\partial}{\partial x} F_0(\mu^\pm(x), t) = \pm \sqrt{\rho(x)} \frac{\partial}{\partial \xi} F_0(\xi, t) \Big|_{\xi=\mu^\pm(x)}. \quad (12)$$

Using the main equation (7) it is easy to prove that

$$A(x, 0) = 0, \quad (13)$$

$$\frac{\sqrt{\rho(x)} - 1}{\sqrt{\rho(x)} + 1} \frac{d}{dx} A(x, \mu^+(x)) = \frac{d}{dx} \{A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)\}. \quad (14)$$

Lemma 1. *The relations*

$$-s''(x, \lambda) + q(x) s(x, \lambda) = \lambda^2 \rho(x) s(x, \lambda), \quad (15)$$

$$s(0, \lambda) = 0, \quad s'(0, \lambda) = 1 \quad (16)$$

hold.

Proof. Assume that $b(x) \in W_2^2(0, \pi)$. Differentiating twice on x the identity

$$\begin{aligned} J(x, t) := & \frac{2}{1 + \sqrt{\rho(t)}} A(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A(x, 2a-t) + \\ & + F(x, t) + \int_0^{\mu^+(x)} A(x, \xi) F_0(\xi, t) d\xi = 0, \end{aligned} \quad (17)$$

using (11) we have

$$\begin{aligned} J'_x(x, t) = & \frac{2}{1 + \sqrt{\rho(t)}} A'_x(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A'_x(x, 2a-t) + \\ & + F'_x(x, t) + \int_0^{\mu^+(x)} A'_x(x, \xi) F_0(\xi, t) d\xi + \sqrt{\rho(x)} A(x, \mu^+(x)) F_0(\mu^+(x), t) + \\ & + \sqrt{\rho(x)} F_0(\mu^-(x), t) (A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)) = 0, \\ J''_{xx}(x, t) = & \frac{2}{1 + \sqrt{\rho(t)}} A''_{xx}(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A''_{xx}(x, 2a-t) + \\ & + F''_{xx}(x, t) + \int_0^{\mu^+(x)} A''_{xx}(x, \xi) F_0(\xi, t) d\xi + \sqrt{\rho(x)} F_0(\mu^+(x), t) \frac{\partial A(x, \xi)}{\partial x} \Big|_{\xi=\mu^+(x)} + \\ & + \sqrt{\rho(x)} F_0(\mu^-(x), t) \left(\frac{\partial A(x, \xi)}{\partial x} \Big|_{\xi=\mu^-(x)+0} - \frac{\partial A(x, \xi)}{\partial x} \Big|_{\xi=\mu^-(x)-0} \right) + \\ & + \sqrt{\rho(x)} F_0(\mu^+(x), t) \frac{d}{dx} A(x, \mu^+(x)) + \sqrt{\rho(x)} A(x, \mu^+(x)) \frac{\partial}{\partial x} F_0(\mu^+(x), t) + \\ & + \sqrt{\rho(x)} (A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)) \frac{\partial}{\partial x} F_0(\mu^-(x), t) + \\ & + \sqrt{\rho(x)} F_0(\mu^-(x), t) \frac{d}{dx} \{A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)\} = 0. \end{aligned} \quad (18)$$

Differentiating the identity (17) on t we have

$$\begin{aligned} J'_t(x, t) &= \frac{2\sqrt{\rho(t)}}{1 + \sqrt{\rho(t)}} A'_t(x, \mu^+(t)) - \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A'_t(x, 2a-t) + \\ &\quad + F'_t(x, t) + \int_0^{\mu^+(x)} A(x, \xi) F'_{0t}(\xi, t) d\xi = 0, \\ J''_{tt}(x, t) &= \frac{2\rho(t)}{1 + \sqrt{\rho(t)}} A''_{tt}(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A''_{tt}(x, 2a-t) + \\ &\quad + F''_{tt}(x, t) + \int_0^{\mu^+(x)} A(x, \xi) F''_{0tt}(\xi, t) d\xi = 0. \end{aligned}$$

According to (10)

$$\begin{aligned} J''_{tt}(x, t) &= \frac{2}{1 + \sqrt{\rho(t)}} A''_{tt}(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A''_{tt}(x, 2a-t) + \\ &\quad + \frac{1}{\rho(t)} F''_{tt}(x, t) + \int_0^{\mu^+(x)} A(x, \xi) F''_{0\xi\xi}(\xi, t) d\xi = 0. \end{aligned}$$

Integrating twice by parts, using (11) and (13) we get

$$\begin{aligned} &\int_0^{\mu^+(x)} A(x, \xi) F''_{0\xi\xi}(\xi, t) d\xi = \\ &= (A(x, \mu^-(x) - 0) - A(x, \mu^-(x) + 0)) \frac{\partial}{\partial \xi} F_0(\xi, t) \Big|_{\xi=\mu^-(x)} + \\ &+ A(x, \mu^+(x)) \frac{\partial}{\partial \xi} F_0(\xi, t) \Big|_{\xi=\mu^+(x)} - F_0(x, \mu^-(x)) \frac{\partial A(x, \xi)}{\partial \xi} \Big|_{\xi=\mu^-(x)-0} + \\ &\quad + F_0(x, 0) \frac{\partial A(x, \xi)}{\partial \xi} \Big|_{\xi=0} - F_0(x, \mu^+(x)) \frac{\partial A(x, \xi)}{\partial \xi} \Big|_{\xi=\mu^+(x)} + \\ &\quad + F_0(x, \mu^-(x)) \frac{\partial A(x, \xi)}{\partial \xi} \Big|_{\xi=\mu^-(x)+0} + \int_0^{\mu^+(x)} A''_{\xi\xi}(x, \xi) F_0(\xi, t) d\xi. \end{aligned}$$

Thus we obtain:

$$\begin{aligned} J''_{tt}(x, t) &= \frac{2}{1 + \sqrt{\rho(t)}} A''_{tt}(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A''_{tt}(x, 2a-t) + \\ &+ \frac{1}{\rho(t)} F''_{tt}(x, t) + (A(x, \mu^-(x) - 0) - A(x, \mu^-(x) + 0)) \frac{\partial}{\partial \xi} F_0(\xi, t) \Big|_{\xi=\mu^-(x)} + \\ &\quad + A(x, \mu^+(x)) \frac{\partial}{\partial \xi} F_0(\xi, t) \Big|_{\xi=\mu^+(x)} - F_0(x, \mu^-(x)) \frac{\partial A(x, \xi)}{\partial \xi} \Big|_{\xi=\mu^-(x)-0} + \\ &\quad + F_0(x, 0) \frac{\partial A(x, \xi)}{\partial \xi} \Big|_{\xi=0} - F_0(x, \mu^+(x)) \frac{\partial A(x, \xi)}{\partial \xi} \Big|_{\xi=\mu^+(x)} + \end{aligned}$$

$$+F_0(x, \mu^-(x)) \left. \frac{\partial A(x, \xi)}{\partial \xi} \right|_{\xi=\mu^-(x)+0} + \int_0^{\mu^+(x)} A''_{\xi\xi}(x, \xi) F_0(\xi, t) d\xi. \quad (19)$$

Using (17) – (19) in the identity

$$J''_{xx}(x, t) - \rho(x) J''_{tt}(x, t) - q(x) J(x, t) \equiv 0$$

we have

$$\begin{aligned} & \frac{2}{1 + \sqrt{\rho(t)}} A''_{xx}(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A''_{xx}(x, 2a-t) + F''_{xx}(x, t) + \\ & + \int_0^{\mu^+(x)} A''_{xx}(x, \xi) F_0(\xi, t) d\xi + \sqrt{\rho(x)} F_0(\mu^+(x), t) \left. \frac{\partial A(x, \xi)}{\partial x} \right|_{\xi=\mu^+(x)} + \\ & + \sqrt{\rho(x)} F_0(\mu^-(x), t) \left(\left. \frac{\partial A(x, \xi)}{\partial x} \right|_{\xi=\mu^-(x)+0} - \left. \frac{\partial A(x, \xi)}{\partial x} \right|_{\xi=\mu^-(x)-0} \right) + \\ & + \sqrt{\rho(x)} F_0(\mu^+(x), t) \frac{d}{dx} A(x, \mu^+(x)) + \sqrt{\rho(x)} A(x, \mu^+(x)) \frac{\partial}{\partial x} F_0(\mu^+(x), t) + \\ & + \sqrt{\rho(x)} (A(x, \mu^-(x)+0) - A(x, \mu^-(x)-0)) \frac{\partial}{\partial x} F_0(\mu^-(x), t) + \\ & + \sqrt{\rho(x)} F_0(\mu^-(x), t) \frac{d}{dx} \{A(x, \mu^-(x)+0) - A(x, \mu^-(x)-0)\} - \\ & - \rho(x) \left(\frac{2}{1 + \sqrt{\rho(t)}} A''_{tt}(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A''_{tt}(x, 2a-t) + \right. \\ & + \frac{1}{\rho(t)} F''_{tt}(x, t) + (A(x, \mu^-(x)-0) - A(x, \mu^-(x)+0)) \left. \frac{\partial}{\partial \xi} F_0(\xi, t) \right|_{\xi=\mu^-(x)} + \\ & + A(x, \mu^+(x)) \left. \frac{\partial}{\partial \xi} F_0(\xi, t) \right|_{\xi=\mu^+(x)} - F_0(x, \mu^-(x)) \left. \frac{\partial A(x, \xi)}{\partial \xi} \right|_{\xi=\mu^-(x)-0} + \\ & + F_0(x, 0) \left. \frac{\partial A(x, \xi)}{\partial \xi} \right|_{\xi=0} - F_0(x, \mu^+(x)) \left. \frac{\partial A(x, \xi)}{\partial \xi} \right|_{\xi=\mu^+(x)} + \\ & + F_0(x, \mu^-(x)) \left. \frac{\partial A(x, \xi)}{\partial \xi} \right|_{\xi=\mu^-(x)+0} + \int_0^{\mu^+(x)} A''_{\xi\xi}(x, \xi) F_0(\xi, t) d\xi \Big) - \\ & - q(x) \left(\frac{2}{1 + \sqrt{\rho(t)}} A(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A(x, 2a-t) + \right. \\ & \left. + F(x, t) + \int_0^{\mu^+(x)} A(x, \xi) F_0(\xi, t) d\xi \right) = 0, \end{aligned}$$

from which using the formulas (6), (9) – (12) and (14) we get

$$\begin{aligned} & \left[\frac{2}{1 + \sqrt{\rho(t)}} A''_{xx}(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A''_{xx}(x, 2a-t) - \right. \\ & -\rho(x) \left(\frac{2}{1 + \sqrt{\rho(t)}} A''_{tt}(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A''_{tt}(x, 2a-t) \right) - \\ & \left. -q(x) \left(\frac{2}{1 + \sqrt{\rho(t)}} A(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a-t)}}{1 + \sqrt{\rho(2a-t)}} A(x, 2a-t) \right) \right] + \\ & + \int_0^{\mu^+(x)} [A''_{xx}(x, \xi) - \rho(x) A''_{\xi\xi}(x, \xi) - q(x) A(x, \xi)] F_0(\xi, t) d\xi = 0. \end{aligned} \quad (20)$$

Homogeneous equation (20) has only trivial solution (see [2]), i.e.

$$A''_{xx}(x, t) - \rho(x) A''_{tt}(x, t) - q(x) A(x, t) = 0, \quad 0 < t < x. \quad (21)$$

Differentiating (8) twice we calculate

$$\begin{aligned} s'(x, \lambda) = & s'_0(x, \lambda) + \int_0^{\mu^+(x)} A'_x(x, t) \frac{\sin \lambda t}{\lambda} dt + \sqrt{\rho(x)} A(x, \mu^+(x)) \frac{\sin \lambda \mu^+(x)}{\lambda} + \\ & + \frac{\sqrt{\rho(x)} \sin \lambda \mu^-(x)}{\lambda} \{A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)\}, \end{aligned} \quad (22)$$

$$\begin{aligned} s''(x, \lambda) = & s''_0(x, \lambda) + \int_0^{\mu^+(x)} A''_{xx}(x, t) \frac{\sin \lambda t}{\lambda} dt + \sqrt{\rho(x)} \frac{\partial}{\partial x} A(x, t) \Big|_{t=\mu^+(x)} \frac{\sin \lambda \mu^+(x)}{\lambda} + \\ & + \sqrt{\rho(x)} \left(\frac{\partial}{\partial x} A(x, t) \Big|_{t=\mu^-(x)+0} - \frac{\partial}{\partial x} A(x, t) \Big|_{t=\mu^-(x)-0} \right) \frac{\sin \lambda \mu^-(x)}{\lambda} + \\ & + \frac{\sqrt{\rho(x)} \sin \lambda \mu^+(x)}{\lambda} \frac{d}{dx} A(x, \mu^+(x)) + \rho(x) A(x, \mu^+(x)) \cos \lambda \mu^+(x) + \\ & + \frac{\sqrt{\rho(x)} \sin \lambda \mu^-(x)}{\lambda} \frac{d}{dx} \{A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)\} - \\ & - \rho(x) \{A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)\} \cos \lambda \mu^-(x). \end{aligned} \quad (23)$$

On the other hand, integrating twice by parts, using (13) we have

$$\begin{aligned} \lambda^2 \rho(x) s(x, \lambda) = & \lambda^2 \rho(x) s_0(x, \lambda) + \lambda^2 \rho(x) \int_0^{\mu^+(x)} A(x, t) \frac{\sin \lambda t}{\lambda} dt = \\ = & -s''_0(x, \lambda) + \rho(x) \cos \lambda \mu^-(x) [A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)] - \end{aligned}$$

$$\begin{aligned}
 & -\rho(x) \cos \lambda \mu^+(x) A(x, \mu^+(x)) + \rho(x) \frac{\sin \lambda \mu^-(x)}{\lambda} \left[\frac{\partial}{\partial t} A(x, t) \Big|_{t=\mu^-(x)-0} - \right. \\
 & \left. - \frac{\partial}{\partial t} A(x, t) \Big|_{t=\mu^-(x)+0} \right] + \rho(x) \frac{\sin \lambda \mu^+(x)}{\lambda} \frac{\partial}{\partial t} A(x, t) \Big|_{t=\mu^+(x)} - \\
 & -\rho(x) \int_0^{\mu^+(x)} A''_{tt}(x, t) \frac{\sin \lambda t}{\lambda} dt. \tag{24}
 \end{aligned}$$

According to (8), (23) – (24) we get

$$\begin{aligned}
 s''(x, \lambda) + \lambda^2 \rho(x) s(x, \lambda) - q(x) s(x, \lambda) = & s''_0(x, \lambda) + \int_0^{\mu^+(x)} A''_{xx}(x, t) \frac{\sin \lambda t}{\lambda} dt + \\
 & + \sqrt{\rho(x)} \frac{\partial}{\partial x} A(x, t) \Big|_{t=\mu^+(x)} \frac{\sin \lambda \mu^+(x)}{\lambda} + \\
 & + \sqrt{\rho(x)} \left(\frac{\partial}{\partial x} A(x, t) \Big|_{t=\mu^-(x)+0} - \frac{\partial}{\partial x} A(x, t) \Big|_{t=\mu^-(x)-0} \right) \frac{\sin \lambda \mu^-(x)}{\lambda} + \\
 & + \frac{\sqrt{\rho(x)} \sin \lambda \mu^+(x)}{\lambda} \frac{d}{dx} A(x, \mu^+(x)) + \rho(x) A(x, \mu^+(x)) \cos \lambda \mu^+(x) + \\
 & + \frac{\sqrt{\rho(x)} \sin \lambda \mu^-(x)}{\lambda} \frac{d}{dx} \{A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)\} - \\
 -\rho(x) \{ & A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0) \} \cos \lambda \mu^-(x) - s''_0(x, \lambda) + \rho(x) \cos \lambda \mu^-(x) \times \\
 \times [& A(x, \mu^-(x) + 0) - A(x, \mu^-(x) - 0)] - \rho(x) \cos \lambda \mu^+(x) A(x, \mu^+(x)) + \\
 & + \rho(x) \frac{\sin \lambda \mu^-(x)}{\lambda} \left[\frac{\partial}{\partial t} A(x, t) \Big|_{t=\mu^-(x)-0} - \frac{\partial}{\partial t} A(x, t) \Big|_{t=\mu^-(x)+0} \right] + \\
 & + \rho(x) \frac{\sin \lambda \mu^+(x)}{\lambda} \frac{\partial}{\partial t} A(x, t) \Big|_{t=\mu^+(x)} - \rho(x) \int_0^{\mu^+(x)} A''_{tt}(x, t) \frac{\sin \lambda t}{\lambda} dt - \\
 & -q(x) \left[\frac{1}{2} \left(1 + \frac{1}{\sqrt{\rho(x)}} \right) \frac{\sin \lambda \mu^+(x)}{\lambda} + \right. \\
 & \left. + \frac{1}{2} \left(1 - \frac{1}{\sqrt{\rho(x)}} \right) \frac{\sin \lambda \mu^-(x)}{\lambda} + \int_0^{\mu^+(x)} A(x, t) \frac{\sin \lambda t}{\lambda} dt \right].
 \end{aligned}$$

Hence using (9), (14) and (21) we obtain (15). For $x = 0$ formulas (8) and (22) imply (16). Lemma 1 is proved in the case $b(x) \in W_2^2(0, \pi)$.

The proof of Lemma 1 in the case $b(x) \in W_2^1(0, \pi)$ is carried out by a standard method (see e.g. [3], p. 56).

Lemma 2. For each function $g(x) \in L_2(0, \pi; \rho)$ the relation

$$\int_0^\pi \rho(x) g^2(x) dx = \sum_{n=1}^\infty \frac{1}{\alpha_n} \left(\int_0^\pi \rho(t) g(t) s(t, \lambda_n) dt \right)^2 \quad (25)$$

holds.

Proof. Using formula (5),

$$F(x, t) = \sum_{n=1}^\infty \left(\frac{s_0(t, \lambda_n) s_0(x, \lambda_n)}{\alpha_n} - \frac{s_0(t, \lambda_n^0) s_0(x, \lambda_n^0)}{\alpha_n^0} \right), \quad (26)$$

and

$$\frac{\sin \lambda \xi}{\lambda} = \begin{cases} s_0(\xi, \lambda), & \xi < a, \\ \frac{2\alpha}{1+\alpha} s_0\left(\frac{\xi}{\alpha} + a - \frac{a}{\alpha}, \lambda\right) + \frac{1-\alpha}{1+\alpha} s_0(2a - \xi, \lambda), & \xi > a \end{cases}$$

it is easy to transform formula (8) to the form

$$s(x, \lambda) = s_0(x, \lambda) + \int_0^x \rho(t) \Phi(x, t) s_0(t, \lambda) dt, \quad (27)$$

and the main equation (7) to the form

$$\Phi(x, t) + F(x, t) + \int_0^x \rho(\xi) \Phi(x, \xi) F(\xi, t) d\xi = 0, \quad (28)$$

where

$$\Phi(x, t) = \frac{2}{1 + \sqrt{\rho(t)}} A(x, \mu^+(t)) + \frac{1 - \sqrt{\rho(2a - t)}}{1 + \sqrt{\rho(2a - t)}} A(x, 2a - t).$$

Solving relation (27) with respect to $s_0(x, \lambda)$ we get

$$s_0(x, \lambda) = s(x, \lambda) + \int_0^x \rho(t) H(x, t) s(t, \lambda) dt. \quad (29)$$

By the standard method we can prove that

$$H(x, t) = F(x, t) + \int_0^t \rho(\xi) \Phi(t, \xi) F(x, \xi) d\xi, \quad 0 \leq t \leq x, \quad (30)$$

where $H(x, t)$ is defined by (29).

Denote $Q(\lambda) = \int_0^\pi \rho(t) g(t) s(t, \lambda) dt$ and using (27) transform it to the form

$$Q(\lambda) = \int_0^\pi \rho(t) h(t) s_0(t, \lambda) dt,$$

where

$$h(t) = g(t) + \int_t^\pi \rho(\xi) g(\xi) \Phi(\xi, t) d\xi. \quad (31)$$

In a similar way, using (29) we get

$$g(t) = h(t) + \int_t^\pi \rho(\xi) h(\xi) H(\xi, t) d\xi. \quad (32)$$

According to (31) we have

$$\begin{aligned} \int_0^\pi \rho(t) h(t) F(x, t) dt &= \int_0^\pi \left[g(t) + \int_t^\pi \rho(\xi) g(\xi) \Phi(\xi, t) d\xi \right] \rho(t) F(x, t) dt = \\ &= \int_0^\pi \rho(t) g(t) \left[F(x, t) + \int_0^t \rho(\xi) \Phi(t, \xi) F(x, \xi) d\xi \right] dt = \\ &= \int_0^x \rho(t) g(t) \left[F(x, t) + \int_0^t \rho(\xi) \Phi(t, \xi) F(x, \xi) d\xi \right] dt + \\ &+ \int_x^\pi \rho(t) g(t) \left[F(x, t) + \int_0^t \rho(\xi) \Phi(t, \xi) F(x, \xi) d\xi \right] dt. \end{aligned}$$

Using (28) and (30) we get

$$\int_0^\pi \rho(t) h(t) F(x, t) dt = \int_0^x \rho(t) g(t) H(x, t) dt - \int_x^\pi \rho(t) g(t) \Phi(t, x) dt. \quad (33)$$

From (26) and Parseval's equality we obtain

$$\begin{aligned} &\int_0^\pi \rho(t) h^2(t) dt + \int_0^\pi \int_0^\pi \rho(t) \rho(x) h(t) h(x) F(x, t) dx dt = \\ &= \int_0^\pi \rho(t) h^2(t) dt + \sum_{n=1}^\infty \left[\frac{1}{\alpha_n} \left(\int_0^\pi \rho(t) h(t) s_0(t, \lambda_n) dt \right)^2 - \right. \\ &\quad \left. - \frac{1}{\alpha_n^0} \left(\int_0^\pi \rho(t) h(t) s_0(t, \lambda_n^0) dt \right)^2 \right] = \\ &= \sum_{n=1}^\infty \frac{1}{\alpha_n} \left(\int_0^\pi \rho(t) h(t) s_0(t, \lambda_n) dt \right)^2 = \sum_{n=1}^\infty \frac{Q^2(\lambda_n)}{\alpha_n}. \end{aligned}$$

According to (33)

$$\begin{aligned} \sum_{n=1}^{\infty} \frac{Q^2(\lambda_n)}{\alpha_n} &= \int_0^{\pi} \rho(t) h^2(t) dt + \int_0^{\pi} \rho(x) h(x) \left[\int_0^x \rho(t) g(t) H(x,t) dt \right] dx - \\ &- \int_0^{\pi} \rho(x) h(x) \left[\int_x^{\pi} \rho(t) g(t) \Phi(t,x) dt \right] dx = \int_0^{\pi} \rho(t) h^2(t) dt + \\ &+ \int_0^{\pi} \rho(t) g(t) \left[\int_t^{\pi} \rho(x) h(x) H(x,t) dx \right] dt - \\ &- \int_0^{\pi} \rho(x) h(x) \left[\int_x^{\pi} \rho(t) g(t) \Phi(t,x) dt \right] dx, \end{aligned}$$

whence by the formulas (31), (32)

$$\begin{aligned} \sum_{n=1}^{\infty} \frac{Q^2(\lambda_n)}{\alpha_n} &= \int_0^{\pi} \rho(t) h^2(t) dt + \int_0^{\pi} \rho(t) g(t) (g(t) - h(t)) dt - \\ &- \int_0^{\pi} \rho(x) h(x) (h(x) - g(x)) dx = \int_0^{\pi} \rho(t) g^2(t) dt, \end{aligned}$$

i.e. relation (25) holds.

Lemma 2 is proved.

Corollary. For any functions $f(x), g(x) \in L_2(0, \pi; \rho)$ the relation

$$\int_0^{\pi} \rho(x) f(x) g(x) dx = \sum_{n=1}^{\infty} \frac{1}{\alpha_n} \int_0^{\pi} \rho(t) f(t) s(t, \lambda_n) dt \int_0^{\pi} \rho(t) g(t) s(t, \lambda_n) dt \quad (34)$$

holds.

Lemma 3. The relation

$$\int_0^{\pi} \rho(t) s(t, \lambda_n) s(t, \lambda_k) dt = \begin{cases} 0, & n \neq k, \\ \alpha_n, & n = k \end{cases} \quad (35)$$

holds.

Proof. 1) Let $f(x) \in W_2^2[0, \pi]$, $f(0) = f(\pi) = 0$. Consider the series

$$f^*(x) = \sum_{n=1}^{\infty} c_n s(x, \lambda_n), \quad (36)$$

where

$$c_n = \frac{1}{\alpha_n} \int_0^{\pi} \rho(x) f(x) s(x, \lambda_n) dx. \quad (37)$$

Using Lemma 1 and integrating twice by parts, we get

$$\begin{aligned} c_n &= \frac{1}{\alpha_n \lambda_n^2} \int_0^\pi f(x) (-s''(x, \lambda_n) + q(x) s(x, \lambda_n)) dx = \\ &= \frac{1}{\alpha_n \lambda_n^2} \left[f'(\pi) s(\pi, \lambda_n) + \int_0^\pi s(x, \lambda_n) (-f''(x) + f(x) q(x)) dx \right]. \end{aligned}$$

Applying asymptotic formulas (4) and

$$s(x, \lambda) = O\left(\frac{1}{|\lambda|} e^{|\operatorname{Im} \lambda| \mu^+(x)}\right), \quad s'(x, \lambda) = O\left(e^{|\operatorname{Im} \lambda| \mu^+(x)}\right)$$

we can see that as $n \rightarrow \infty$

$$c_n = O\left(\frac{1}{n^3}\right), \quad s(x, \lambda_n) = O\left(\frac{1}{n}\right)$$

uniformly on $[0, \pi]$. Hence, series (36) converges absolutely and uniformly on $[0, \pi]$. By (34) and (37) we have

$$\begin{aligned} \int_0^\pi \rho(x) f(x) g(x) dx &= \sum_{n=1}^\infty c_n \int_0^\pi \rho(t) g(t) s(t, \lambda_n) dt = \\ &= \int_0^\pi \rho(t) g(t) \sum_{n=1}^\infty c_n s(t, \lambda_n) dt = \int_0^\pi \rho(t) g(t) f^*(t) dt. \end{aligned}$$

By arbitrariness of $g(x)$ we conclude that $f^*(x) = f(x)$, i.e.

$$f(x) = \sum_{n=1}^\infty c_n s(x, \lambda_n). \tag{38}$$

2) Fix $k \geq 1$ and assume $f(x) = s(x, \lambda_k)$. Then, by (38)

$$s(x, \lambda_k) = \sum_{n=1}^\infty c_{nk} s(x, \lambda_n),$$

where

$$c_{nk} = \frac{1}{\alpha_n} \int_0^\pi \rho(x) s(x, \lambda_n) s(x, \lambda_k) dx.$$

We can show that the system of functions $\{s_0(x, \lambda_n)\}_{n \geq 1}$ is minimal in $L_2(0, \pi; \rho)$. Then by (8) the system of functions $\{s(x, \lambda_n)\}_{n \geq 1}$ is minimal in $L_2(0, \pi; \rho)$ as well. Therefore $c_{nk} = \delta_{nk}$ (δ_{nk} is a Kronecker symbol) and we arrive at (35).

Lemma 3 is proved.

Lemma 4. For all $n \geq 1$ the equality

$$s(\pi, \lambda_n) = 0$$

holds.

Proof. It is easy to show that

$$(\lambda_n^2 - \lambda_m^2) \int_0^\pi \rho(x) s(x, \lambda_n) s(x, \lambda_m) dx = s(\pi, \lambda_n) s'(\pi, \lambda_m) - s'(\pi, \lambda_n) s(\pi, \lambda_m).$$

By (35) we get

$$s(\pi, \lambda_n) s'(\pi, \lambda_m) - s'(\pi, \lambda_n) s(\pi, \lambda_m) = 0. \quad (39)$$

We shall prove that for any n , $s'(\pi, \lambda_n) \neq 0$. Assume the contrary, i.e. there exists such m that $s'(\pi, \lambda_m) = 0$. Then from (39) we have $s'(\pi, \lambda_n) = 0$ for $n \neq m$.

On the other hand

$$s'(\pi, \lambda_n) = s'_0(\pi, \lambda_n) + \int_0^{\mu^+(\pi)} A'_x(\pi, t) \frac{\sin \lambda_n t}{\lambda_n} dt + \frac{\alpha \sin \lambda_n \mu^+(\pi)}{\lambda_n} A(\pi, \mu^+(\pi)) + \\ + \frac{\alpha \sin \lambda_n \mu^-(\pi)}{\lambda_n} \{A(\pi, \mu^-(\pi) + 0) - A(\pi, \mu^-(\pi) - 0)\},$$

i.e. $s'(\pi, \lambda_n) \sim s'_0(\pi, \lambda_n^0) \neq 0$ as $n \rightarrow \infty$, that contradicts the condition $s'(\pi, \lambda_n) = 0$, $n \neq m$. Thus $s'(\pi, \lambda_n) \neq 0$ for any n , and from (39) we have

$$\frac{s(\pi, \lambda_n)}{s'(\pi, \lambda_n)} = \frac{s(\pi, \lambda_m)}{s'(\pi, \lambda_m)} \stackrel{def}{=} H,$$

i.e. for any n , $s'(\pi, \lambda_n) H = s(\pi, \lambda_n)$. Since $s(\pi, \lambda_n) = o(1)$ as $n \rightarrow \infty$, we have $H = 0$, i.e. $s(\pi, \lambda_n) = 0$.

Lemma 4 is proved.

It follows from Lemmas 1, 3 and 4 that the numbers $\{\lambda_n^2, \alpha_n\}_{n \geq 1}$ are spectral data of the constructed boundary value problem $L(q(x))$.

The theorem is proved.

Algorithm of the construction of the function $q(x)$ by spectral data $\{\lambda_n^2, \alpha_n\}$ follows from the proof of the theorem:

- 1) By the given numbers $\{\lambda_n^2, \alpha_n\}_{n \geq 1}$ the functions $F_0(x, t)$ and $F(x, t)$ are constructed by the formulas (5) and (6), respectively;
- 2) The function $A(x, t)$ is found from equation (7);
- 3) $q(x)$ is calculated by the formula (9).

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