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**THE SOLUTION OF A NONLOCAL BOUNDARY
VALUE PROBLEM FOR A FOURTH ORDER
PARTIAL DIFFERENTIAL EQUATION**

Abstract

In the paper we prove existence of a classic solution of a non-local boundary value problem for a fourth order partial differential equation.

For the equation [1]

$$u_{tt}(x, t) - u_{ttxx}(x, t) - u_{xx}(x, t) = f(x, t) \tag{1}$$

in the domain $D_T = \{(x, t) : 0 \leq x \leq 1, 0 \leq t \leq T\}$ we consider a problem under the conditions

$$\begin{aligned} u(x, 0) + \delta u(x, T) &= \varphi(x), \\ u_t(x, 0) + \delta u_t(x, T) &= \psi(x), \quad 0 \leq x \leq 1 \end{aligned} \tag{2}$$

and boundary conditions

$$u(x, 0) = 0, \quad u_x(1, t) + du_{xx}(1, t) = 0, \quad 0 \leq t \leq T, \tag{3}$$

where $d > 0$, δ are the given numbers $f(x, t)$, $\varphi(x)$, $\psi(x)$ are the given functions, $u(x, t)$ is the desired function. Under the classic solution of the problem (1) – (3) we understand the function $u(x, t)$ continuous in a closed domain D_T together with all derivatives entering in to the equation (1) and satisfying all the conditions (1) – (3) in the ordinary sense.

Let's consider a spectral problem [2]

$$y''(x) + \lambda y(x) = 0, \quad 0 \leq x \leq 1, \tag{4}$$

$$y(0) = 0, \quad y'(1) = d\lambda y(1), \quad d > 0 \tag{5}$$

that has only ligen functions

$$y_k(x) = \sqrt{2} \sin(\sqrt{\lambda_k}x), \quad k = 0, 1, 2, \dots,$$

with positive eigen numbers λ_k from the equation $ctg\sqrt{\lambda} = d\sqrt{\lambda}$. We give a zero index to any eigen function and enumerate the other ones in increase order of eigen numbers.

The following statements are formulated and justified in the paper [2].

Lemma 1. *Beginning with some number N it holds the estimate*

$$0 < \sqrt{\lambda_k} - k\pi < (dk\pi)^{-1}.$$

Lemma 2. *The system $\{z_k(x)\}$ conjugated to the system*

$$\{y_k(x)\}, \quad k = 1, 2, \dots,$$

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is determined by the formula

$$z_k(x) = \sqrt{2} \left(\sin(\sqrt{\lambda_k}x) - \sin \sqrt{\lambda_k} \sin(\sqrt{\lambda_0}x) / \sin \sqrt{\lambda_0} \right) / \left(1 + d \sin^2 \sqrt{\lambda_n} \right).$$

Theorem 1. The system $\{y_k(x)\}$, $k = 1, 2, \dots$ forms a basis in the space $L_p(0, 1)$, $p > 1$, for $p = 2$ even the Riesz basis.

In the paper [3], it is known that

$$\sqrt{\lambda_k} = k\pi + O\left(\frac{1}{k}\right),$$

$$y_k(x) = \sqrt{2} \sin(k\pi x) + O\left(\frac{1}{k}\right),$$

$$\sqrt{2} \cos \sqrt{\lambda_k}x = \sqrt{2} \cos(k\pi x) + O\left(\frac{1}{k}\right).$$

Assume $g(x) \in L_p(0, 1)$. Then

$$\begin{aligned} & \left(\sum_{k=1}^{\infty} \left(\sqrt{2} \int_0^1 g(x) \sin \sqrt{\lambda_k}x dx \right)^2 \right)^{1/2} \leq \\ & \leq \left(\frac{\sqrt{2}}{2} + \left(\sum_{k=1}^{\infty} O(k^{-2}) \right)^{1/2} \right) \|g(x)\|_{L_2(0,1)}, \end{aligned} \quad (6)$$

$$\begin{aligned} & \left(\sum_{k=1}^{\infty} \left(\sqrt{2} \int_0^1 g(x) \cos \sqrt{\lambda_k}x dx \right)^2 \right)^{1/2} \leq \\ & \leq \left(\frac{\sqrt{2}}{2} + \left(\sum_{k=1}^{\infty} O(k^{-2}) \right)^{1/2} \right) \|g(x)\|_{L_2(0,1)}. \end{aligned} \quad (7)$$

Theorem 2. [4]. Let the function $g(x)$ satisfy the conditions

$$g(x) \in W_2^{2n}(0, 1), \quad g(0) = g'(0) = \dots = g^{(2n-2)}(0) = 0, \quad n \geq 1,$$

$$J \equiv g(1) + \left(1/d \sin \sqrt{\lambda_0} \right) \int_0^1 g(x) \sin \sqrt{\lambda_0}x dx = 0,$$

$$g''(1) + g'(1)/d = 0, \dots, g^{(2n-2)}(1) + g^{(2n-3)}(1)/d = 0, \quad n \geq 2.$$

then the Fourier series

$$\sum_{k=1}^{\infty} g_k y_k(x)$$

converges in the metric $W_2^{2n}(0, 1)$.

But if the function $g(x)$ is subjected to requirements:

$$g(x) \in W_2^{(2n+1)}(0,1), \quad g(0) = g''(0) = \dots = g^{(2n)}(0) = 0, \quad n \geq 0, \quad J = 0,$$

$$g''(0) + g'(1)/d = 0, \dots, g^{(2n)}(1) + g^{(2n-1)}(1)/d = 0, \quad n \geq 1,$$

then the Fourier series

$$\sum_{k=1}^{\infty} g_k y_k(x)$$

converges in the metric $W_2^{(2n+1)}(0,1)$ where

$$g_k = \int_0^1 g(x) z_k(x) dx.$$

Let $g(x) \in L_2(0,1)$, then

$$\begin{aligned} g_k &= \frac{\sqrt{2}}{1 + d \sin^2 \sqrt{\lambda_k}} \left(\int_0^1 g(x) \sin \sqrt{\lambda_k} x dx - \frac{\sin \sqrt{\lambda_k}}{\sin \sqrt{\lambda_0}} \int_0^1 g(x) \sin \sqrt{\lambda_0} x dx \right) = \\ &= \frac{\sqrt{2}}{1 + d \sin^2 \sqrt{\lambda_k}} \left(\int_0^1 g(x) \sin \sqrt{\lambda_k} x dx - \frac{\cos \sqrt{\lambda_k}}{d \sqrt{\lambda_k} \sin \sqrt{\lambda_0}} \int_0^1 g(x) \sin \sqrt{\lambda_0} x dx \right). \end{aligned}$$

Hence, allowing for (6) we have:

$$\left(\sum_{k=1}^{\infty} g_k^2 \right)^{1/2} \leq M_0 \|g(x)\|_{L_2(0,1)}, \quad (8)$$

where

$$M_0 = \left(\frac{\sqrt{2}}{2} + \left(\sum_{k=1}^{\infty} 0(k^{-2}) \right)^{1/2} + \frac{1}{d \sqrt{6} \sin \sqrt{\lambda_0}} \right) \sqrt{2}.$$

Assume that $g'(x) \in L_2(0,1)$ and

$$g(0) = 0, \quad J = 0. \quad (9)$$

It is known that [3],

$$g_k = \frac{\sqrt{2}}{1 + d \sin^2 \sqrt{\lambda_k}} \frac{1}{\sqrt{\lambda_k}} \int_0^1 g'(x) \cos \sqrt{\lambda_k} x dx. \quad (10)$$

Hence, allowing for (7) we find:

$$\left(\sum_{k=1}^{\infty} \left(\sqrt{\lambda_k} |g_k| \right)^2 \right)^{1/2} \leq M_1 \|g'(x)\|_{L_2(0,1)}, \quad (11)$$

where

$$M_1 = \frac{\sqrt{2}}{2} + \left(\sum_{k=1}^{\infty} 0(k^{-2}) \right)^{1/2}.$$

We assume that for $g(x) \in C^{-1}[0, 1]$, $g''(x) \in L_2(0, 1)$ the conditions (9) are fulfilled and obtain the equality

$$g_k = \frac{-\sqrt{2}}{1 + d \sin^2 \sqrt{\lambda_k}} \frac{1}{\lambda_k} \int_0^1 g''(x) \sin \sqrt{\lambda_k} x dx + \frac{\cos \sqrt{\lambda_k}}{d \sqrt{\lambda_k} \lambda_k} g'(1).$$

Hence, allowing for (6) we have:

$$\left(\sum_{k=1}^{\infty} (\lambda_k |g_k|)^2 \right)^{1/2} \leq M_2 \|g''(x)\|_{L_2(0,1)} + \frac{1}{d\sqrt{6}} |g'(1)|. \quad (12)$$

Now, assume that $g(x) \in C^2[0, 1]$, $g'''(x) \in L_2(0, 1)$ and

$$g(0) = g''(0) = 0, \quad J = 0, \quad g''(1) + g'(1)/d = 0.$$

Then

$$g_k = -\frac{\sqrt{2}}{1 + d \sin^2 \sqrt{\lambda_k}} \frac{1}{\lambda_k \sqrt{\lambda_k}} \int_0^1 g'''(x) \cos \sqrt{\lambda_k} x dx. \quad (13)$$

Hence, allowing for (7) we find:

$$\left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} |g_k|)^2 \right)^{1/2} \leq M_1 \|g'''(x)\|_{L_2(0,1)}. \quad (14)$$

We'll look for the classic solution of the problem (1) – (2) in the form

$$u(t, x) = \sum_{k=1}^{\infty} u_k(t) y_k(x), \quad (15)$$

where

$$u_k(t) = \int_0^1 u(x, t) z_k(x) dx \quad (16)$$

is a solution of the following problem:

$$(1 + \lambda_k) u_k''(t) + \lambda_k u_k(t) = f_k(t) \quad (17)$$

$$u_k(0) + \delta u_k(T) = \varphi_k, \quad u_k'(0) + \delta u_k'(T) = \psi_k \quad (k = 1, 2, \dots), \quad (18)$$

where

$$f_k(t) = \int_0^1 f(x, t) z_k(x) dx, \quad \varphi_k = \int_0^1 \varphi(x) z_k(x) dx, \quad \psi_k = \int_0^1 \psi(x) z_k(x) dx.$$

Solving the problem (17) – (18) we find:

$$\begin{aligned}
 u_k(t) = & \frac{1}{\beta_k \rho_k(T)} \{ \beta_k (\cos \beta_k t + \delta \cos \beta_k (T - t)) \varphi_k + \\
 & (\sin \beta_k t - \delta \sin \beta_k (T - t)) \psi_k - \frac{\delta}{1 + \lambda_k} \int_0^T f_k(\tau) (\sin \beta_k (T + t - \tau) + \\
 & + \delta \sin \beta_k (t - \tau)) d\tau \} + \frac{1}{\beta_k (1 + \lambda_k)} \int_0^t f_k(\tau) \sin \beta_k (t - \tau) d\tau \quad (k = 1, 2), \quad (19)
 \end{aligned}$$

where

$$\beta_k = \sqrt{\frac{\lambda_k}{1 + \lambda_k}}, \quad \rho_k(T) = 1 + 2\delta \cos \beta_k T + \delta^2 \quad (k = 1, 2, \dots).$$

Obviously

$$\begin{aligned}
 u'_k(t) = & \frac{1}{\rho_k(T)} \{ \beta_k (-\sin \beta_k t + \delta \sin \beta_k (T - t)) \varphi_k + (\cos \beta_k t - \\
 & - \delta \cos \beta_k (T - t)) \psi_k - \frac{\delta}{1 + \lambda_k} \int_0^T f_k(\tau) (\cos \beta_k (T + t - \tau) + \delta \cos \beta_k (t - \tau)) d\tau \} + \\
 & + \frac{1}{1 + \lambda_k} \int_0^t f_k(\tau) \cos \beta_k (t - \tau) d\tau \quad (k = 1, 2), \quad (20)
 \end{aligned}$$

$$\begin{aligned}
 u''_k(t) = & \frac{1}{1 + \lambda_k} f_k(\tau) - \frac{\beta_k}{\rho_k(T)} \{ \beta_k (\cos \beta_k t + \delta \cos \beta_k (T - t)) \varphi_k + \\
 & + (\sin \beta_k t - \delta \sin \beta_k (T - t)) \psi_k - \frac{\delta}{1 + \lambda_k} \int_0^T f_k(\tau) (\sin \beta_k (T + t - \tau) + \\
 & + \delta \sin \beta_k (t - \tau)) d\tau \} - \frac{\beta_k}{1 + \lambda_k} \int_0^t f_k(\tau) \sin \beta_k (t - \tau) d\tau \quad (k = 1, 2, \dots). \quad (21)
 \end{aligned}$$

The following theorem is true.

Theorem 3. Let

1. $\varphi(x) \in C^2[0, 1]$, $\varphi'''(x) \in L_2(0, 1)$, $\varphi(0) = \varphi''(0) = 0$,
 $\varphi(1) + (1/(d \sin \sqrt{\lambda_0})) \int_0^1 \varphi(x) \sin \sqrt{\lambda_0} x dx = 0$, $\varphi''(1) + \varphi'(1)/d = 0$;
2. $\psi(x) \in C^2[0, 1]$, $\psi'''(x) \in L_2(0, 1)$, $\psi(0) = \psi''(0) = 0$,
 $\psi(1) + (1/(d \sin \sqrt{\lambda_0})) \int_0^1 \psi(x) \sin \sqrt{\lambda_0} x dx = 0$, $\psi''(1) + \psi'(1)/d = 0$;
3. $f(x, t) \in C(D_T)$, $f_x(x, t) \in L_2(D_T)$, $f(0, t) = 0$,
 $f(1, t) + (1/(d \sin \sqrt{\lambda_0})) \int_0^1 f(x, t) \sin \sqrt{\lambda_0} x dx = 0$, $(t \in [0, T])$.

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4. $\delta \neq \pm 1$, $1 + 2\delta \cos \beta_k T + \delta^2 \neq 0$ ($k = 1, 2, \dots$).

Then the function

$$\begin{aligned}
u(x, t) = & \sum_{k=1}^{\infty} \left\{ \frac{1}{\beta_k \rho_k(T)} [\beta_k (\cos \beta_k t + \delta \cos \beta_k (T-t)) \varphi_k] + \right. \\
& + (\sin \beta_k t - \delta \sin \beta_k (T-t)) \psi_k - \\
& \left. - \frac{\delta}{1 + \lambda_k} \int_0^T f_k(\tau) (\sin \beta_k (T+t-\tau) + \delta \sin \beta_k (t-\tau)) d\tau \right\} + \\
& \left. + \frac{1}{\beta_k (1 + \lambda_k)} \int_0^t f_k(\tau) \sin \beta_k (t-\tau) d\tau \right\} y_k(x) \quad (22)
\end{aligned}$$

is a classic solution of the problem (1) – (3).

Proof. Obviously

$$1/\sqrt{2} < \beta_k < 1, \quad |\rho_k(T)| \geq 1 + 2\delta^2 - |\delta| \equiv 1/\rho.$$

Taking these into account, from (19), (20), (21) we have:

$$\begin{aligned}
|u_k(t)| & \leq \sqrt{2}(1 + |\delta|)(|\varphi_k| + |\psi_k|) + \\
& + \sqrt{2T}(1 + \rho|\delta|(1 + |\delta|)) \frac{1}{\lambda_k} \left(\int_0^T |f_k(\tau)|^2 d\tau \right)^{1/2}, \\
|u'_k(t)| & \leq \frac{1}{\lambda_k} |f_k(t)| + \rho(1 + |\delta|)(|\varphi_k| + |\psi_k|) + \\
& + (1 + \rho|\delta|(1 + |\delta|)) \frac{\sqrt{T}}{\lambda_k} \left(\int_0^T |f_k(\tau)|^2 d\tau \right)^{1/2}.
\end{aligned}$$

Hence we get

$$\begin{aligned}
& \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} \|u_k(t)\|_{C[0,T]})^2 \right)^{1/2} \leq \sqrt{6}\rho(1 + |\delta|) \times \\
& \times \left[\left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} |\varphi_k|)^2 \right)^{1/2} + \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} |\psi_k|)^2 \right)^{1/2} \right] + \\
& + \sqrt{6}(1 + \rho|\delta|(1 + |\delta|)) \sqrt{T} \left(\int_0^T \sum_{k=1}^{\infty} (\lambda_k^{1/2} |f_k(\tau)|)^2 d\tau \right)^{1/2}, \quad (23) \\
& \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} \|u'_k(t)\|_{C[0,T]})^2 \right)^{1/2} \leq \sqrt{3}\rho(1 + |\delta|) \times
\end{aligned}$$

$$\begin{aligned} & \times \left[\left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} |\varphi_k|)^2 \right)^{1/2} + \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} |\psi_k|)^2 \right)^{1/2} \right] + \\ & \sqrt{3} (1 + \rho |\delta| (1 + |\delta|)) \sqrt{T} \left(\int_0^T \sum_{k=1}^{\infty} (\lambda_k^{1/2} |f_k(\tau)|)^2 d\tau \right)^{1/2}, \end{aligned} \quad (24)$$

$$\begin{aligned} & \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} \|u_k''(t)\|_{C[0,T]})^2 \right)^{1/2} \leq \sqrt{2} \left(\sum_{k=1}^{\infty} (\lambda_k^{1/2} |f_k(t)|)^2 \right)^{1/2} + 2\rho (1 + |\delta|) \times \\ & \times \left[\left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} |\varphi_k|)^2 \right)^{1/2} + \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} |\psi_k|)^2 \right)^{1/2} \right] + \\ & + 2 (1 + \rho |\delta| (1 + |\delta|)) \sqrt{T} \left(\int_0^T \sum_{k=1}^{\infty} (\lambda_k^{1/2} |f_k(\tau)|)^2 d\tau \right)^{1/2}. \end{aligned} \quad (25)$$

Then from (23) – (25) allowing for (11), (14) we find:

$$\begin{aligned} & \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} \|u_k(t)\|_{C[0,T]})^2 \right)^{1/2} \leq \\ & \leq \sqrt{6}\rho (1 + |\delta|) M_1 \left(\|\varphi'''(x)\|_{L_2(0,1)} + \|\psi'''(x)\|_{L_2(0,1)} \right) + \\ & + \sqrt{6T} (1 + \rho |\delta| (1 + |\delta|)) M_1 \|f_x(x, t)\|_{L_2(D_T)}, \end{aligned} \quad (26)$$

$$\begin{aligned} & \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} \|u_k'(t)\|_{C[0,T]})^2 \right)^{1/2} \leq \\ & \leq \sqrt{3}\rho (1 + |\delta|) M_1 \left(\|\varphi'''(x)\|_{L_2(0,1)} + \|\psi'''(x)\|_{L_2(0,1)} \right) + \\ & + \sqrt{3T} (1 + \rho |\delta| (1 + |\delta|)) M_1 \|f_x(x, t)\|_{L_2(D_T)}, \end{aligned} \quad (27)$$

$$\begin{aligned} & \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} \|u_k''(t)\|_{C[0,T]})^2 \right)^{1/2} \leq 2M_1 \left\| \|f_x(x, t)\|_{L_2(0,1)} \right\|_{C[0,T]} + \\ & + 2\rho (1 + |\delta|) M_1 \left(\|\varphi'''(x)\|_{L_2(0,1)} + \|\psi'''(x)\|_{L_2(0,1)} \right) + \\ & + 2 (1 + \rho |\delta| (1 + |\delta|)) \sqrt{T} M_1 \|f_x(x, t)\|_{L_2(D_T)}. \end{aligned} \quad (28)$$

Obviously

$$|u(x, t)| \leq \left(\sum_{k=1}^{\infty} \lambda_k^{-3} \right)^{1/2} \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} \|u_k(t)\|_{C[0,T]})^2 \right)^{1/2}, \quad (29)$$

$$|u_t(x, t)| \leq \left(\sum_{k=1}^{\infty} \lambda_k^{-3} \right)^{1/2} \left(\sum_{k=1}^{\infty} (\lambda_k^{3/2} \|u_k'(t)\|_{C[0,T]})^2 \right)^{1/2}, \quad (30)$$

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$$|u_{xx}(x, t)| \leq \left(\sum_{k=1}^{\infty} \lambda_k^{-1} \right)^{1/2} \left(\sum_{k=1}^{\infty} \left(\lambda_k^{3/2} \|u_k(t)\|_{C[0,T]} \right)^2 \right)^{1/2}, \quad (31)$$

$$|u_{tt}(x, t)| \leq \left(\sum_{k=1}^{\infty} \lambda_k^{-3} \right)^{1/2} \left(\sum_{k=1}^{\infty} \left(\lambda_k^{3/2} \|u_k''(t)\|_{C[0,T]} \right)^2 \right)^{1/2}, \quad (32)$$

$$|u_{ttxx}(x, t)| \leq \left(\sum_{k=1}^{\infty} \lambda_k^{-1} \right)^{1/2} \left(\sum_{k=1}^{\infty} \left(\lambda_k^{3/2} \|u_k''(t)\|_{C[0,T]} \right)^2 \right)^{1/2}. \quad (33)$$

From (29) – (33) allowing for (26) – (28) it follows that the functions $u(x, t)$, $u_t(x, t)$, $u_{xx}(x, t)$, $u_{tt}(x, t)$, $u_{ttxx}(x, t)$ are continuous in D_T . We can see by direct verification that the functions $u(x, t)$ satisfy the equation and conditions (2), (3) in the ordinary sense. The theorem is proved.

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