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DEFINITION OF THE UNKNOWN COEFFICIENT OF A PARABOLIC EQUATION WITH NON-LOCAL BOUNDARY AND COMPLEMENTARY CONDITIONS

Abstract

The inverse non-selfadjoint boundary value problem is investigated for linear parabolic equations.

Consider on the domain $Q = (0,1) \times (0,T)$ the problem

$$\begin{cases} u_t(t,x) - u_{xx}(t,x) + a(t)u(t,x) = F(t,x), & (1) \end{cases}$$

$$\begin{cases} u(0,x) = \varphi(x), \quad 0 \leq x \leq 1, & (2) \end{cases}$$

$$\begin{cases} u(t,0) = 0, \quad u_x(t,0) = u_x(t,1), \quad 0 \leq t \leq T, & (3) \end{cases}$$

$$\begin{cases} u_x(t,0) + \beta u(t,1) = g(t), \quad 0 \leq t \leq T, & (4) \end{cases}$$

where $0 < T < +\infty$; F, φ, g are given functions, β is a given number, and $u(t,x), a(t)$ are desired functions ([1], [2], [3], [4], [5]).

Definition. Under the classic solution of the problem (1)-(4) we understand the pair $\{u(t,x), a(t)\}$ of functions $u(t,x), a(t)$ possessing the properties:

a) $u(t,x), u_t(t,x), u_{xx}(t,x) \in C(\overline{Q})$;

b) $a(t) \in C([0, T])$;

c) all conditions of (1)-(4) are satisfied in a classic sense.

To investigate the problem (1)-(4) we must study the following spectral problem

$$\begin{cases} X''(x) + \lambda X(x) = 0, \quad 0 \leq x \leq 1, & (5) \end{cases}$$

$$\begin{cases} X(0) = 0, \quad X'(1) = X'(0). & (6) \end{cases}$$

It is clear that the problem (5), (6) is non-selfadjoint. Corresponding conjugate problem has the form

$$\begin{cases} Y''(x) + \lambda Y(x) = 0, & (7) \end{cases}$$

$$\begin{cases} Y(0) = Y(1), \quad Y'(1) = 0. & (8) \end{cases}$$

It is obvious that the problem (5), (6) has eigen values

$$\lambda_k = (2\pi k)^2, \quad k = 0, 1, 2, \dots$$

and eigen-functions

$$\overline{X}_0(x) = x, \quad \overline{X}_k(x) = \sin 2\pi kx, \quad k = 1, 2, \dots \quad (9)$$

A system of eigen functions $\overline{X}_k(x), (k = 0, 1, 2, \dots)$ doesn't form a basis in $L_2(0,1)$ ([2]). Complement the eigen functions $\overline{X}_k(x), (k = 0, 1, 2, \dots)$ with respect to the whole system by the adjoined functions

$$\overline{\overline{X}}_k(x) = x \cos 2\pi kx, \quad k = 1, 2, \dots \quad (10)$$

of the boundary value problem (5), (6).

Over denote the functions systems (9) and (10) as follows

$$X_0(x) = x, \quad X_{2k-1}(x) = x \cos 2\pi kx, \quad X_{2k}(x) = \sin 2\pi kx, \quad k = 1, 2, \dots \quad (11)$$

Further, we find eigen and adjoint functions of the problem (7), (8) and over denote as follows:

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$$Y_0(x) = 2, \quad Y_{2k-1}(x) = 4 \cos 2\pi kx, \quad Y_{2k}(x) = 4(1-x) \sin 2\pi kx, \quad (12)$$

$$k = 1, 2, \dots$$

It is easy to calculate that sequences (11) and (12) form a biorthogonal system of functions on the interval (0,1), i.e.

$$(X_i, Y_j) \equiv \int_0^1 X_i(x) Y_j(x) dx = \delta_{ij}. \quad (13)$$

Here δ_{ij} is a Kronecker's symbol.

Thus, the system (11) forms a basis in $L_2(0,1)$ and the system (12) forms a biorthogonal system of functions in $L_2(0,1)$, and it is obvious that for each classic solution $\{u(t,x), a(t)\}$ of the problem (1)-(4) its first component $u(t,x)$ has the form

$$u(t,x) = \sum_{k=0}^{\infty} u_k(t) X_k(x), \quad (14)$$

where

$$u_k(t) = \int_0^1 u(t,x) Y_k(x) dx, \quad k = 0, 1, \dots \quad (15)$$

By multiplying both sides of the equality (1) by $Y_k(x)$ ($k = 0, 1, 2, \dots$) and by integrating with respect to x from zero to the unit we obtain the following equation system with respect to $u_k(t)$ ($k = 0, 1, 2, \dots$):

$$u'_0(t) + a(t)u_0(t) = F_0(t), \quad (16)$$

$$u'_{2k-1}(t) + [(2\pi k)^2 + a(t)]u_{2k-1}(t) = F_{2k-1}(t), \quad (17)$$

$$u'_{2k}(t) + [(2\pi k)^2 + a(t)]u_{2k}(t) + 4\pi k u_{2k-1}(t) = F_{2k}(t), \quad k = 1, 2, \dots, \quad (18)$$

where

$$F_k(t) = \int_0^1 F(t,x) Y_k(x) dx, \quad k = 0, 1, 2, \dots \quad (19)$$

By solving the equation system (16), (17) and (18) we get

$$u_0(t) = A_0 e^{-\int_0^t a(s) ds} + \int_0^t F_0(\tau) e^{-\int_0^{\tau} a(s) ds} d\tau, \quad (20)$$

$$u_{2k-1}(t) = A_{2k-1} e^{-(2\pi k)^2 t - \int_0^t a(s) ds} + \int_0^t F_{2k-1}(\tau) e^{-(2\pi k)^2 (t-\tau) - \int_0^{\tau} a(s) ds} d\tau, \quad (21)$$

$$u_{2k}(t) = [A_{2k} - 4\pi k A_{2k-1} t] e^{-(2\pi k)^2 t - \int_0^t a(s) ds} + \int_0^t [F_{2k}(\tau) - 4\pi k F_{2k-1}(\tau)(t-\tau)] e^{-(2\pi k)^2 (t-\tau) - \int_0^{\tau} a(s) ds} d\tau, \quad k = 1, 2, \dots \quad (22)$$

where A_k ($k = 0, 1, 2, \dots$) are arbitrary constants. By using the representation (11) and initial condition (2) we get

$$u(0,x) = \sum_{k=0}^{\infty} u_k(0) X_k(x) = \sum_{k=0}^{\infty} \varphi_k X_k(x) = \varphi(x), \quad (23)$$

where $\varphi_k = \int_0^1 \varphi(x) Y_k(x) dx$.

It follows from (23) that

$$u_k(0) = \varphi_k, \quad k = 0, 1, 2, \dots \tag{24}$$

Using the initial condition (24), appearing in (20)-(22) determine the constants A_k ($k = 0, 1, 2, \dots$) in the form

$$A_k = \varphi_k, \quad k = 0, 1, 2, \dots \tag{25}$$

It follows from (14), (20)-(22) and (25) that

$$\begin{aligned} u(t, x) = & \left[\varphi_0 e^{-\int_0^t a(s) ds} + \int_0^t F_0(\tau) e^{-\int_\tau^t a(s) ds} d\tau \right] X_0(x) + \\ & + \sum_{k=1}^{\infty} \left[\varphi_{2k-1} e^{-(2\pi k)^2 t - \int_0^t a(s) ds} + \int_0^t F_{2k-1}(\tau) e^{-(2\pi k)^2 (t-\tau) - \int_\tau^t a(s) ds} d\tau \right] \times \\ & \times X_{2k-1}(x) + \sum_{k=1}^{\infty} \left\{ [\varphi_{2k} - 4\pi k \varphi_{2k-1} t] e^{-(2\pi k)^2 t - \int_0^t a(s) ds} + \right. \\ & \left. + \int_0^t [F_{2k}(\tau) - 4\pi k (t-\tau) F_{2k-1}(\tau)] e^{-(2\pi k)^2 (t-\tau) - \int_\tau^t a(s) ds} d\tau \right\} X_{2k}(x). \end{aligned} \tag{26}$$

To determine the second component $a(t)$ of the solution, by using (4) we find

$$u_{xt}(t, 0) + \beta u_t(t, 1) = g'(t).$$

Hence we shall have

$$\begin{aligned} g'(t) + a(t)g(t) = & (\beta + 1)F_0(t) + \sum_{k=1}^{\infty} [2\pi k F_{2k}(t) + (\beta + 1)(2\pi k)^2 F_{2k-1}(t)] - \\ & - \sum_{k=1}^{\infty} [(2\pi k)^3 \varphi_{2k} + (\beta + 3)(2\pi k)^2 \varphi_{2k-1} - 2(2\pi k)^4 t \varphi_{2k-1}] \cdot e^{-(2\pi k)^2 t - \int_0^t a(s) ds} - \\ & - \sum_{k=1}^{\infty} \int_0^t [(2\pi k)^3 F_{2k}(\tau) + (\beta + 3)(2\pi k)^2 F_{2k-1}(\tau) - 2(2\pi k)^4 (t-\tau) F_{2k-1}(\tau)] \times \\ & \times e^{-(2\pi k)^2 (t-\tau) - \int_\tau^t a(s) ds} d\tau. \end{aligned} \tag{27}$$

It is valid the following

Theorem. Let

- 1) $\varphi(x) \in C^V([0, 1])$, $\varphi(0) = \varphi''(0) = \varphi^{IV}(0) = 0$, $\varphi'(0) = \varphi'(1)$, $\varphi'''(0) = \varphi'''(1)$;
- 2) $g(t) \in C'([0, T])$, $g(t) \neq 0 \quad \forall t \in [0, T]$;
- 3) $F(x, t) \in C(\overline{Q})$ and at each fixed $t \in [0, T]$
 $F(\cdot, x) \in C^V([0, 1])$, $F(\cdot, 0) = F_{xx}(\cdot, 0) = F_{xxx}(\cdot, 0) = 0$, $F_x(\cdot, 0) = F_x(\cdot, 1)$,
 $F_{xx}(\cdot, 0) = F_{xx}(\cdot, 1)$.

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Then under sufficiently small values of T the problem (1)-(4) has a classic solution.

Proof. Write (27) in the form of the operator equation

$$a(t) = P[a(t)]. \quad (28)$$

It is obvious that for any $a(t)$ from the space $C([0, T])$, $P[a(t)] \in C([0, T])$, i.e.

$$P: C([0, T]) \rightarrow C([0, T]).$$

Now prove the compressibility of the operator P .

Let $a(t), b(t) \in C([0, T])$.

Then

$$\begin{aligned} |P[a(t)] - P[b(t)]| &\leq \frac{1}{|g(t)|} \sum_{k=1}^{\infty} \left\{ (2\pi k)^3 |\varphi_{2k}| + ((\beta + 3)(2\pi k)^2 + 2T(2\pi k)^4) |\varphi_{2k-1}| \right\} \times \\ &\times \left| e^{-\int_0^t a(s) ds} - e^{-\int_0^t b(s) ds} \right| + \frac{1}{|g(t)|} \int_0^t \left\{ (2\pi k)^3 |F_{2k}(\tau)| + ((\beta + 3)(2\pi k)^2 + 4T(2\pi k)^4) \times \right. \\ &\times \left. |F_{2k-1}(\tau)| \right\} \left| e^{-\int_{\tau}^t a(s) ds} - e^{-\int_{\tau}^t b(s) ds} \right| d\tau. \end{aligned}$$

By c_i ($i = \overline{1, 5}$) denote positive constants independent on t .

By virtue of conditions 1)-3)

$$\begin{aligned} \sum_{k=1}^{\infty} \left\{ (2\pi k)^3 |\varphi_{2k}| + ((\beta + 3)(2\pi k)^2 + 2T(2\pi k)^4) |\varphi_{2k-1}| \right\} &\leq c_1, \\ \int_0^T \left\{ \sum_{k=1}^{\infty} \left[(2\pi k)^3 |F_{2k}(\tau)| + ((\beta + 3)(2\pi k)^2 + 4T(2\pi k)^4) |F_{2k-1}(\tau)| \right] \right\} d\tau &\leq c_2, \end{aligned}$$

$$\max_{0 \leq t \leq T} \frac{1}{|g(t)|} = c_3,$$

$$\left| e^{-\int_0^t a(s) ds} - e^{-\int_0^t b(s) ds} \right| \leq c_4 \cdot T \max_{0 \leq t \leq T} |a(t) - b(t)|,$$

$$\left| e^{-\int_{\tau}^t a(s) ds} - e^{-\int_{\tau}^t b(s) ds} \right| \leq c_5 \cdot T \max_{0 \leq t \leq T} |a(t) - b(t)|.$$

Then

$$\max_{0 \leq t \leq T} |P[a(t)] - P[b(t)]| \leq c_3 (c_1 c_4 + c_2 c_5) T \max_{0 \leq t \leq T} |a(t) - b(t)|.$$

Choosing T sufficiently small, we obtain that

$$c_3 (c_1 c_4 + c_2 c_5) T = q < 1.$$

Consequently, $P[a(t)]$ is a compressible operator. Then by a fixed-point theorem, the equation (28) has a unique solution $a(t)$ from $C([0, T])$.

Substituting the found solution (27) in (26) we find the first component $u(t, x)$ of the solution $\{u(t, x), a(t)\}$. It is easy to prove that the pair $\{u(t, x), a(t)\}$ satisfies (1)-(4).

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