

Ahmad H. JAMSHIDIPOUR, Hidayat M. HUSEYNOV

SOLUTION OF THE INVERSE SCATTERING PROBLEM FOR THE STURM-LIOUVILLE EQUATION WITH A SPECTRAL PARAMETER IN DISCONTINUITY CONDITION

Abstract

In the paper, a total solution of the inverse scattering problem is given for Sturm-Liouville equation with a spectral parameter in the discontinuity condition in the absence of a discrete spectrum.

It is known that the inverse scattering problem for the differential equation

$$-y'' - \lambda p(x)y + q(x)y = \lambda^2 y, \quad -\infty < x < \infty, \quad (1)$$

where $p(x)$ and $q(x)$ are real functions, $p(x)$ is absolutely continuous, and $p(x)$, $p'(x)$, $q(x)$ rather rapidly decrease as $|\lambda| \rightarrow \infty$, was investigated in the papers [1]-[3]. In the case $p(x) \equiv 0$, under the condition

$$\int_{-\infty}^{\infty} (1 + |x|) |q(x)| dx < \infty \quad (2)$$

the papers [4]-[6] deal with the solution of the inverse scattering problem.

In the present paper, we consider the inverse scattering problem for equation (1), when $p(x) = \beta\delta(x - a)$, where β is a real number, $\delta(x - a)$ is a Dirac function, and condition (2) is fulfilled. In this case, equation (1) may be reduced to the problem ([7])

$$-y'' + q(x)y = \lambda^2 y, \quad -\infty < x < \infty \quad (3)$$

$$y(a + 0) = y(a - 0), \quad (4)$$

$$y'(a + 0) - y'(a - 0) = \lambda\beta y(a). \quad (5)$$

As is known [8], problem (3)-(5) for all λ from the half-plane $\text{Im } \lambda \geq 0$ has the solutions $e^+(x, \lambda)$, $e^-(x, \lambda)$ representable in the form

$$e^\pm(x, \lambda) = e_0^\pm(x, \lambda) \pm \int_x^{\pm\infty} K^\pm(x, t) e^{\pm i\lambda t} dt, \quad (6)_\pm$$

moreover, the kernels $K^\pm(x, t)$ satisfy the inequalities

$$|K^\pm(x, t)| \leq \frac{C}{2} \sigma^\pm \left(\frac{x+t}{2} \right) e^{C\sigma_{1^\pm}(x)}, \quad 0 < |x - a| < \pm(t - a), \quad (7)_\pm$$

$$|K^\pm(x, t)| \leq \left\{ \frac{C}{2} \sigma^\pm \left(\frac{x+t}{2} \right) + \frac{|\beta|}{4} \sigma^\pm \left(\frac{x+2a-t}{2} \right) \right\} e^{C\sigma_{1^\pm}(x)}, \quad |t - a| < \pm(a - x),$$

where $C = 1 + \frac{|\beta|}{2}$, $\sigma^\pm(x) = \pm \int_x^{\pm\infty} |q(\xi)| d\xi$, $\sigma_1^\pm(x) = \pm \int_x^{\pm\infty} \sigma^\pm(\xi) d\xi$.

Furthermore, the functions $K^\pm(x, t)$ are continuous for $t \neq 2a - x$, $x \neq a$, and the following relations are fulfilled:

$$K^\pm(x, t) = \pm \frac{1}{2} \int_x^{\pm\infty} q(\xi) d\xi, \quad \pm x > \pm a,$$

$$K^\pm(x, x) = \pm \frac{1}{2} \left(1 + \frac{i\beta}{2}\right) \int_x^{\pm\infty} q(\xi) d\xi, \quad \pm x < \pm a, \quad (8)_\pm$$

$$K^\pm(x, 2a - x + 0) - K^\pm(x, 2a - x - 0) =$$

$$= \mp \frac{i\beta}{2} \left\{ \int_x^a q(\xi) d\xi - \int_x^{\pm\infty} q(\xi) d\xi \right\}, \quad \pm x < \pm a.$$

In formulae (6)_±, the functions $e_0^\pm(x, \lambda)$ are the Jost solutions of problem (3)-(5) for $q(x) \equiv 0$:

$$e_0^\pm(x, \lambda) = \begin{cases} e^{\pm i\lambda x}, & \pm x > \pm a \\ \left(1 + \frac{i\beta}{2}\right) e^{\pm i\lambda x} - \frac{i\beta}{2} e^{\pm i\lambda(2a-x)}, & \pm x < \pm a. \end{cases}$$

For $\lambda \in R/\{0\}$ the following relations hold:

$$\frac{1}{a(\lambda)} e^\pm(x, \lambda) = r^\mp(\lambda) e^\mp(x, \lambda) + \overline{e^\mp(x, \lambda)}, \quad (9)_\pm$$

where

$$r^-(\lambda) = \frac{b(\lambda)}{a(\lambda)}, \quad r^+(\lambda) = -\frac{\overline{b(\lambda)}}{a(\lambda)},$$

$$a(\lambda) = \frac{1}{2i\lambda} W[e^+(x, \lambda), e^-(x, \lambda)], \quad b(\lambda) = -\frac{1}{2i\lambda} W[e^+(x, \lambda), \overline{e^-(x, \lambda)}],$$

$W[y_1, y_2] = y_1'y_2 - y_1'y_2'$ is the wronskian of the functions y_1 and y_2 .

The functions $r^-(\lambda)$, $r^+(\lambda)$ and $\frac{1}{a(\lambda)}$ are called left and right reflection factor and passage factor, respectively. The function $a(\lambda)$ admits analytic continuation to the half-plane $\text{Im } \lambda > 0$ and may have there at most finite number of zeros. In future, we'll assume that the zeros also are absent, i.e.

$$a(\lambda) \neq 0 \quad \text{for} \quad \text{Im } \lambda > 0, \quad (10)$$

and consequently, the discrete spectrum of problem (3)-(5) is absent.

In this case, the inverse scattering problem for problem (3)-(5) is in renewal of the function $q(x)$ by the left or right reflection factor and finding necessary and sufficient conditions to which should satisfy an arbitrarily taken function $r(\lambda)$ for it to be the right (left) reflection factor of some problem of the form (3)-(5) with a real coefficient satisfying condition (2).

In the following theorem we give the solution of the inverse scattering problem.

Theorem. For the function $r^+(\lambda)$, $-\infty < \lambda < \infty$ to be the right reflection factor of the problem of the form (3)-(5), without a discrete spectrum, with a real potential $q(x)$ satisfying inequality (2) and with a real number β , it is necessary and sufficient that the following conditions to be fulfilled:

1) for real $\lambda \neq 0$, the function $r^+(\lambda)$ is continuous,

$$|r^+(\lambda)| \leq 1 - c\lambda^2 (1 + 2^2)^{-1} \quad \text{and} \quad r^+(\lambda) - r_0^+(\lambda) = O\left(\frac{1}{\lambda}\right) \quad \text{as } \lambda \rightarrow \pm\infty,$$

where

$$r_0^+(\lambda) = -\frac{i\beta}{2 + i\beta} e^{-2i\lambda a};$$

2) the function $za(z)$, where

$$a(z) = \left(1 + \frac{i\beta}{2}\right) e^{-\frac{1}{2\pi i} \int_{-\infty}^{\pm\infty} \frac{1 - |r^+(\lambda)|^2}{\lambda - z} d\lambda}$$

is continuous in the closed upper half-plane, and

$$\lim_{\lambda \rightarrow 0} \lambda a(\lambda) [r^+(\lambda) + 1] = 0;$$

3) the functions

$$R^\pm(x) = \frac{1}{2\pi} \int_{-\infty}^{\pm\infty} [r^\pm(\lambda) - r_0^\pm(\lambda)] e^{\pm i\lambda x} d\lambda,$$

where $r^-(\lambda) = \overline{-r^+(\lambda) \cdot \frac{\bar{a}(\lambda)}{a(\lambda)}}$, are absolutely continuous on any segment not containing the point $2a$, the derivatives $R^{+\prime}(\lambda)$ and $R^{-\prime}(\lambda)$ for all $\alpha > -\infty$ and $\alpha' < \infty$ satisfy the inequality

$$\int_{\alpha}^{\infty} (1 + |x|) |R^{+\prime}(x)| dx < \infty, \quad \int_{-\infty}^{\alpha'} (1 + |x|) |R^{-\prime}(x)| dx < \infty;$$

4) the solutions $K^\pm(x, y)$ of the equations

$$R_1^\pm(x, y) + \overline{K^\pm(x, y)} \mp \frac{i\beta}{2 + i\beta} K^\pm(x, 2a - y) \pm \int_x^{\pm\infty} K^\pm(x, t) R^\pm(t + y) dt = 0, \quad \pm y > \pm x, \quad (11)_\pm$$

where

$$R_1^\pm(x, y) = \begin{cases} R^\pm(x, y), & \pm x > \pm a \\ \left(1 + \frac{i\beta}{2}\right) R^\pm(x + y) - \frac{i\beta}{2} R^\pm(2a - x + y), & \pm x < \pm a \end{cases},$$

satisfy the conditions

$$K^\pm(x, x)|_{a \mp 0} = \left(1 + \frac{i\beta}{2}\right) K^\pm(x + y)|_{a \mp 0}.$$

Proof. Necessity follows from the results of the papers [8]-[9]. Give the proof of sufficiency.

1. From conditions 1), 3) of the theorem it follows that the main equations (11)₊ and (11)₋ constructed by the function $r^+(\lambda)$, according to the paper [10] have unique solutions $K^+(x, y)$ and $K^-(x, y)$. In the expanded form the main equation (11)_± looks as follows:

$$R^\pm(x + y) + \overline{K^\pm(x, y)} \pm \int_x^{\pm\infty} K^\pm(x, t) R^\pm(t + y) dt = 0, \quad \pm x > \pm a, \quad \pm y > \pm x \quad (12)_\pm$$

$$\begin{aligned} & \left(1 + \frac{i\beta}{2}\right) R^\pm(x + y) - \frac{i\beta}{2} R^\pm(2a - x + y) + \overline{K^\pm(x, y)} \pm \\ & \pm \int_x^{\pm\infty} K^\pm(x, t) R^\pm(t + y) dt = 0, \quad \pm x < \pm a, \quad \pm y > \pm(2a - x). \end{aligned} \quad (13)_\pm$$

$$\begin{aligned} & \left(1 + \frac{i\beta}{2}\right) R^\pm(x + y) - \frac{i\beta}{2} R^\pm(2a - x + y) + \overline{K^\pm(x, y)} \mp \frac{i\beta}{2 + i\beta} K^\pm(x, 2a - y) \pm \\ & \pm \int_x^{\pm\infty} K^\pm(x, t) R^\pm(t + y) dt = 0, \quad \pm x < \pm a, \quad \pm x < y < \pm(2a - x). \end{aligned} \quad (14)_\pm$$

Assuming $y = 2a - x \mp 0$ and $y = 2a - x \pm 0$ in main equations (14)_± and (13)_±, respectively, we have

$$\begin{aligned} & \left(1 + \frac{i\beta}{2}\right) R^\pm(2a \mp 0) - \frac{i\beta}{2} R^\pm(4a - 2x \mp 0) + \overline{K^\pm(x, 2a - x \mp 0)} \mp \\ & \mp \frac{i\beta}{2 + i\beta} K^\pm(x, x \pm 0) \pm \int_x^{\pm\infty} K^\pm(x, t) R^\pm(t + 2a - x \mp 0) dt = 0, \quad \pm x < \pm a, \\ & \left(1 + \frac{i\beta}{2}\right) R^\pm(2a \pm 0) - \frac{i\beta}{2} R^\pm(4a - 2x \pm 0) + \overline{K^\pm(x, 2a - x \pm 0)} \pm \\ & \pm \int_x^{\pm\infty} K^\pm(x, t) R^\pm(t + 2a - x \pm 0) dt = 0, \quad \pm x < \pm a. \end{aligned}$$

We subtract the second relation from the first one. As a result we get :

$$\begin{aligned} & \left(1 + \frac{i\beta}{2}\right) [R^\pm(2a \mp 0) - R^\pm(2a \pm 0)] + \\ & + \overline{K^\pm(x, 2a - x \mp 0) - K^\pm(x, 2a - x \pm 0)} \mp \end{aligned}$$

$$\mp \frac{i\beta}{2+i\beta} K^\pm(x, x \pm 0), \quad \pm x < \pm a.$$

Now assume (see (8)_±)

$$q^\pm(x) = \begin{cases} = \pm \frac{4}{i\beta} \frac{d}{dx} \left[\frac{\mp \frac{4}{2+i\beta} \frac{dK^\pm(x,x)}{dx}}{K^\pm(x, 2a-x+0) - K^\pm(x, 2a-x-0)} \right], & \pm x < \pm a, \\ \mp 2 \frac{dK^\pm(x,x)}{dx}, & \pm x > \pm a. \end{cases} \quad (15)_\pm$$

2. Show that the functions $e^+(x, \lambda)$, $e^-(x, \lambda)$ constructed with the help of $K^+(x, t)$, $K^-(x, t)$ by formula (6)₊ and (6)₋ satisfy the equations

$$-e^{\pm''}(x, \lambda) + q^\pm(x) e^\pm(x, \lambda) = \lambda^2 e^\pm(x, \lambda) \quad (16)_\pm$$

and conditions

$$a^\pm(a+0, \lambda) = e^\pm(a-0, \lambda), \quad (17)_\pm$$

$$e^{\pm'}(a+0, \lambda) - e^{\pm'}(a-0, \lambda) = \lambda \beta e^\pm(a, \lambda), \quad (18)_\pm$$

moreover

$$\int_{x'}^{\infty} (1+|x|) |q^+(x)| dx < \infty, \quad \int_{-\infty}^{x''} (1+|x|) |q^-(x)| dx < \infty. \quad (19)_\pm$$

At first suppose that the functions $R^\pm(x)$ are twice continuously differentiable, and for all $\alpha' > -\infty$, $\beta' < +\infty$

$$\int_{\alpha'}^{\infty} (1+|x|) |R^{\pm''}(x)| dx < +\infty, \quad \int_{-\infty}^{\beta'} (1+|x|) |R^{\pm''}(x)| dx < \infty. \quad (20)$$

Then the solutions $K^\pm(x, y)$ of main equations (11)_± are twice continuously differentiable for $y \neq 2a - x$ and $x \neq a$, and for each x all partial derivatives of first and second order are summable over y .

Consider the domain $\pm x < \pm a$, $\pm x < \pm y < \pm(2a - x)$. Then the main equations (11)_± take the form of (14)_±. Differentiating these equations twice with respect to y and integrating by parts, we get

$$\begin{aligned} & \left(1 + \frac{i\beta}{2}\right) R^{\pm''}(x+y) - \frac{i\beta}{2} R^{\pm''}(2a-x+y) + \overline{K_{yy}^{\pm''}(x,y)} \mp \frac{i\beta}{2+i\beta} K_{yy}^{\pm''}(x, 2a-y) \mp \\ & \mp K^\pm(x, x) R^{\pm'}(x+y) \mp K^\pm(x, t) \Big|_{t=2a-x-0}^{2a-x+0} \cdot R^{\pm'}(2a-x+y) \pm K_t^{\pm'}(x, t) \Big|_{t=x} \times \\ & \times R^\pm(x+y) \pm K_t^{\pm'}(x, t) \Big|_{t=2a-x-0}^{2a-x+0} R^\pm(2a-x+y) \pm \int_x^{\pm\infty} K_{tt}^{\pm''}(x, t) R^\pm(t+y) dt = 0. \end{aligned}$$

Further, differentiating equation (14)_± twice respect to x , we have

$$\left(1 + \frac{i\beta}{2}\right) R^{\pm''}(x+y) - \frac{i\beta}{2} R^{\pm''}(2a-x+y) + \overline{K_{xx}^{\pm''}(x,y)} \mp$$

$$\begin{aligned} & \mp \frac{i\beta}{2+i\beta} K_{xx}^{\pm''}(x, 2a-y) \mp K^{\pm}(x, x) R^{\pm}(x+y) \mp K^{\pm}(x, x) \cdot R^{\pm'}(x+y) \mp \\ & \quad \mp \left[K^{\pm}(x, t) \Big|_{t=2a-x-0}^{2a-x+0} \right] R^{\pm}(2a-x+y) \pm \\ & \quad \pm \left[K^{\pm}(x, t) \Big|_{t=2a-x-0}^{2a-x+0} \right] R^{\pm'}(2a-x+y) \mp K_x^{\pm'}(x, t) \Big|_{t=x} R^{\pm}(x+y) \pm \\ & \quad \pm \left[K_x^{\pm}(x, t) \Big|_{t=2a-x-0}^{2a-x+0} \right] R^{\pm}(2a-x+y) \pm \int_x^{\pm\infty} K_{xx}^{\pm'}(x, t) R^{\pm}(t+y) dt = 0. \end{aligned}$$

Subtracting from the last equality the previous one, we get

$$\begin{aligned} & \overline{K_{xx}^{\pm''}(x, y)} \mp \frac{i\beta}{2+i\beta} K_{xx}^{\pm''}(x, 2a-y) - \overline{K_{yy}^{\pm''}(x, y)} \pm \frac{i\beta}{2+i\beta} K_{yy}^{\pm''}(x, 2a-y) \mp \\ & \quad \mp 2K^{\pm}(x, x) R^{\pm}(x+y) \mp 2 \left[K^{\pm}(x, t) \Big|_{t=2a-x-0}^{2a-x+0} \right]' R^{\pm}(2a-x+y) \pm \\ & \quad \pm \int_x^{\pm\infty} \left(K_{xx}^{\pm''}(x, t) - K_{tt}^{\pm''}(x, t) \right) R^{\pm}(t+y) dt = 0. \end{aligned} \quad (21)_{\pm}$$

From (15)_± and main equation (14)_±

$$\begin{aligned} & \mp 2K_x^{\pm'}(x, x) R^{\pm}(x+y) \mp 2 \left[K^{\pm}(x, t) \Big|_{t=2a-x-0}^{2a-x+0} \right]' R^{\pm}(2a-x+y) = \\ & \quad = \left(1 + \frac{i\beta}{2} \right) q'' + (x) R^{\pm}(x+y) - \frac{i\beta}{2} q^{\pm}(x) R^{\pm}(2a-x+y) = \\ & \quad = q^{\pm(x)} \left[-\overline{K^{\pm}(x, y)} \pm \frac{i\beta}{2} K^{\pm}(x, 2a-y) \mp \int_x^{\pm\infty} K^{\pm}(x, t) R^{\pm}(t+y) dt \right]. \end{aligned} \quad (22)_{\pm}$$

From (21)_± and (22)_± it follows that the functions

$$h_x^{\pm}(y) = K_{xx}^{\pm''}(x, y) - q^{\pm}(x) K^{\pm}(x, y) - K_{yy}^{\pm''}(x, y)$$

satisfy the equations

$$\overline{h_x^{\pm}(y)} \pm \frac{i\beta}{2+i\beta} h_x^{\pm}(2a-y) \pm \int_x^{\pm\infty} h_x^{\pm}(t) R^{\pm}(t+y) dt = 0,$$

i.e. the functions $h_x^{\pm}(t)$ are summable solutions of homogeneous equations corresponding to (14)_±. Behaving in the same way with equations (12)_± and (13)_±, we get that the solutions of main equations (11)_±, according to [10] satisfy the equation

$$K_{xx}^{\pm''}(x, y) - q^{\pm}(x) K^{\pm}(x, y) - K_{yy}^{\pm''}(x, y) = 0. \quad (23)_{\pm}$$

By condition 4) of the theorem, from (15)_± it follows that the functions $K^{\pm}(x, y)$ satisfy relations (8)_±. Under the made suppositions (20) it is easy to show that

$$\lim_{x+y \rightarrow \pm\infty} K_x^{\pm'}(x, y) = \lim_{x+y \rightarrow \pm\infty} K_x^{\pm'}(x, y) = 0. \quad (24)_{\pm}$$

Now, show that the functions $K^\pm(x, y)$ satisfy the conditions

$$K^\pm(a+0, y) = K^\pm(a-0, y), \quad \pm y > \pm a \quad (25)_\pm$$

$$K_x^{\pm'}(a+0, y) - K_x^{\pm'}(a-0, y) = i\beta K_y^{\pm'}(a, y), \quad \pm y > \pm a. \quad (26)_\pm$$

Assume $x = a \pm 0$ and $x = a \mp 0$ in main equations (12) $_\pm$ and (14) $_\pm$, respectively. From the obtained first relation subtract the second one. As a result we get that the differences $K^\pm(a+0, y) - K^\pm(a-0, y)$ are the solutions of homogeneous equations corresponding to equations (11) $_\pm$, for $x = 0$. Therefore, by [10] we get (25) $_\pm$.

Prove that conditions (26) $_\pm$ are also fulfilled.

First of all note that for the solution of main equations, the following relations are valid:

$$\begin{aligned} K^\pm(x, 2a - x \pm 0)|_{a \mp 0} &= K^\pm(x, x)|_{a \pm 0}, \\ K^\pm(x, 2a - x \mp 0)|_{a \mp 0} &= K^\pm(x, x)|_{a \mp 0}. \end{aligned} \quad (27)_\pm$$

Indeed, in equations (12) $_\pm$ assume at first $y = x$, then $x = a \pm 0$, in equations (13) $_\pm$ at first $y = 2a - x \pm 0$, then $x = a \mp 0$. Further, from these equations, subtracting one from another one, by (25) $_\pm$ we get the first equality from (27) $_\pm$. Assuming once $y = 2a - x \mp 0$, $x = a \mp 0$ and $y = x \pm 0$, $x = a \mp 0$ another time in equations (14) $_\pm$ and subtracting one of the obtained equalities from another one it is easy to get the second relation from (27) $_\pm$.

Now differentiate equations (12) $_\pm$ and (13) $_\pm$ with respect to the variable x , and assume $x = a + 0$ and $x = a - 0$, respectively. As a result we have

$$\begin{aligned} &R^{+'}(a+y) + \overline{K_x^{+'}(a+0, y)} - K^+(x, x)|_{a+0} R^+(a+y) + \\ &+ \int_a^{+\infty} K_x^{+'}(a+0, t) R^+(t+y) dt = 0, \\ &\left(1 + \frac{i\beta}{2}\right) R^+(a+y) + \frac{i\beta}{2} R^{+'}(a+y) + \overline{K_x^{+'}(a-0, y)} - K^+(x, x)|_{a-0} R^+(a+y) + \\ &+ [K^+(x, 2a - x + 0) - K^+(x, 2a - x - 0)]|_{a-0} R^+(a+y) + \\ &+ \int_a^{\pm\infty} K_x^{+'}(a-0, t) R^+(t+y) dt = 0. \end{aligned}$$

Subtracting one these equalities from another one, and taking into account (27) $_\pm$, we get

$$\begin{aligned} &-i\beta R^+(a+y) + \overline{K_x^{+'}(a+0, y)} - K_x^+(a-0, y) + \\ &+ \{2K^+(x, x)|_{a-0} - 2K^+(x, x)|_{a+0}\} R^+(a+y) + \\ &+ \int_a^{\pm\infty} [K_x^{+'}(a+0, t) - K_x^{+'}(a-0, t)] R^+(t+y) dt = 0. \end{aligned}$$

Further, assume $x = a$ in equation (12)₊, and differentiate with respect to y . Then integrating by parts, we get

$$R^{+'}(a+y) + \overline{K_y^{+'}(a,y)} - K^+(x,x) \Big|_{a+0} R^+(a+y) + \\ + \int_a^{+\infty} K_t^{+'}(a,t) R^+(t+y) dt = 0.$$

Multiply these equalities by $-i\beta$ and subtract the obtained relation from the equality next to last. As a result, by condition of the theorem we have

$$\overline{K_x^{+'}(a+0,y) - K_x^{+'}(a-0,y) - i\beta K_y^{+'}(a,y)} + \\ + \int_a^{\pm\infty} \left\{ K_x^{+'}(a+0,t) - K_x^{+'}(a-0,t) - i\beta K_t^{+'}(a,y) \right\} R^+(t+y) dt = 0.$$

Hence, according to the uniqueness theorem on the solution of main equation ([10]) we get relation (26)_±. Similaly, proceeding from equations (12)₋ and (13)₋, relation (26)₋ is established.

Thus, by fulfilling conditions (20), the solutions $K^\pm(x,y)$ of main equations (11)_± satisfy equation (23)_± by relations (8)_±, (25)_±, (26)_± conditions (24)_±. Then for solving these problems we get integral equations from [8] which show that the functions $e^\pm(x,\lambda)$ constructed with the help of $K^\pm(x,t)$ by formulae (6)_± satisfy equations (16)_± and conditions (17)_±, (18)_±.

The case when only conditions 3) of the theorem are fulfilled, may be considered by means of the limit passage (see [3], p. 212).

Finally, show that conditions (19) are also fulfilled. Since for $\pm x > \pm a$ main equations (11)_± have the form of (12)_±, i.e. similar to the case $\alpha = 1$ form and conditions 3) of the theorem is the same as in the case $\alpha = 1$, then it is easy show that relations (19) are valid if $x' \geq a$ and $x' \leq a$ (see [3], p. 209). It remains to show that $q^+(x)$ ($q^-(x)$) are summable in the interval (x', a) ((a, x'')) for each $x' > -\infty$ ($x'' < +\infty$). If we use conditions 3) of the theorem, and integrability of partial derivatives $K_x^{\pm'}$, $K_t^{\pm'}$, these facts are easily established by means of the formula (equivalent to equation (14)_±)

$$K^\pm(x,y) = \left(1 + \frac{\beta^2}{4}\right) \left[\overline{\varphi^\pm(x,y)} \mp \frac{i\beta}{2-i\beta} \varphi^\pm(x, 2a-y) \right],$$

where

$$\varphi^\pm(x,y) = - \left(1 + \frac{i\beta}{2}\right) R^\pm(x+y) + \\ + \frac{i\beta}{2} R^\pm(2a-x+y) \mp \int_x^{\pm\infty} K^\pm(x,y) R^\pm(t+y) dt.$$

3. Now for proving theorem, it suffices to show that for real $\lambda \neq 0$ the functions $e^+(x, \lambda)$ and $e^-(x, \lambda)$ are connected with the equalities

$$r^\pm(\lambda) e^\pm(x, \lambda) + \overline{e^\pm(x, \lambda)} = \frac{1}{a(\lambda)} e^\mp(x, \lambda). \quad (28)_\pm$$

Indeed, from (28)_±, by (16)_± it follows that:

$$q^+(x) = q^-(x) \stackrel{def}{=} q(x), \quad -\infty < x < +\infty,$$

and according to (19) we have:

$$\int_{-\infty}^{\infty} (1 + |x|) |q(x)| dx < +\infty.$$

Show that then $r^+(\lambda)$ and $r^-(\lambda)$ are the right and left (respectively) reflection factors of the constructed problem (3)-(5).

Denote the right and left reflection factors of the constructed problem (3)-(5) by $\tilde{r}^+(\lambda)$ and $\tilde{r}^-(\lambda)$, respectively. The functions $e^+(x, \lambda)$ and $e^-(x, \lambda)$ will be the lost solutions of problem (3)-(5). Therefore, by the results of the direct scattering problem we can write the relations

$$\tilde{r}^\pm(\lambda) e^\pm(x, \lambda) + \overline{e^\pm(x, \lambda)} = \frac{1}{\tilde{a}(\lambda)} e^\mp(x, \lambda). \quad (29)_\pm$$

From (28)_± and (29)_± we have

$$a(\lambda) r^+(\lambda) e^+(x, \lambda) + a(\lambda) \overline{e^+(x, \lambda)} = e^-(x, \lambda),$$

$$\tilde{a}(\lambda) \tilde{r}^+(\lambda) \tilde{e}^+(x, \lambda) + \tilde{a}(\lambda) \overline{\tilde{e}^+(x, \lambda)} = e^-(x, \lambda).$$

Subtracting one of these relations from another one, we have

$$\{a(\lambda) r^+(\lambda) - \tilde{a}(\lambda) \tilde{r}^+(\lambda)\} e^+(x, \lambda) + \{a(\lambda) - \tilde{a}(\lambda)\} \overline{e^+(x, \lambda)} = 0.$$

Since $\lambda \neq 0$, the functions $e^+(x, \lambda)$ and $e^-(x, \lambda)$ are linearly independent, then it follows from the last identity that

$$a(\lambda) r^+(\lambda) - \tilde{a}(\lambda) \tilde{r}^+(\lambda) = 0, \quad a(\lambda) - \tilde{a}(\lambda) = 0$$

i.e. $r^+(\lambda) = \tilde{r}^+(\lambda)$, $a(\lambda) = \tilde{a}(\lambda)$, similarly, from relations (28)₋ and (29)₋ we get $r^-(\lambda) = \tilde{r}^-(\lambda)$.

Further, since the function $a(z)$ according to condition 2) of the theorem has no zeros in upper half-plane, then problem (3)-(5) has no discrete spectrum.

4. Now, prove relations (28)_±. Assume

$$\Phi^\pm(x, y) = R_1^\pm(x + y) \pm \int_x^{\pm\infty} K^\pm(x, t) R^\pm(t + y) dt,$$

where

$$R_1(x, y) = \begin{cases} R^\pm(x + y), & \pm x > \pm a, \\ \left(1 + \frac{i\beta}{2}\right) R^\pm(x + y) - \frac{i\beta}{2} R^\pm(2a - x + y), & \pm x < \pm a. \end{cases}$$

Since $R^\pm(y) \in L_2(-\infty, +\infty)$, then for each fixed x $\Phi(x, y) \in L_2(-\infty, +\infty)$ we have

$$\begin{aligned} & \lim_{N \rightarrow +\infty} \int_{-N}^N \Phi^\pm(x, y) e^{\mp\lambda y} dy = \\ & = [r^\pm(\lambda) - r_0^\pm(\lambda)] \left[e_0^\pm(x, \lambda) \pm \int_x^{\pm\infty} K^\pm(x, t) e^{\pm\lambda t} dt \right] = \\ & = [r^\pm(\lambda) - r_0^\pm(\lambda)] e^\pm(x, \lambda). \end{aligned} \tag{30}_\pm$$

On the other hand, from equations (11)_±

$$\Phi^\pm(x, y) = -\overline{K^\pm(x, y)} \pm \frac{i\beta}{2 + i\beta} K^\pm(x, 2a - y), \quad \pm y > \pm x.$$

Therefore

$$\begin{aligned} & \lim_{N \rightarrow +\infty} \int_{-N}^N \Phi^\pm(x, y) e^{\mp\lambda y} dy = \\ & = \lim_{N \rightarrow +\infty} \left\{ \pm \int_{\mp N}^x \Phi^\pm(x, y) e^{\mp\lambda y} dy \right\} - \int_x^{\pm\infty} \overline{K^\pm(x, t)} e^{\mp\lambda y} dy + \\ & + \frac{i\beta}{2 + i\beta} \int_x^{\pm\infty} K^\pm(x, 2a - y) e^{\mp\lambda y} dy = \lim_{N \rightarrow +\infty} \left\{ \pm \int_{\mp N}^x \Phi^\pm(x, y) e^{\mp\lambda y} dy \right\} - \\ & - \overline{e^\pm(x, \lambda)} + \overline{e_0^\pm(x, \lambda)} - r_0^\pm(\lambda) [e^\pm(x, \lambda) + e_0^\pm(x, \lambda)]. \end{aligned} \tag{31}_\pm$$

Comparing (30)_± and (31)_±, and taking into account the formulae

$$r_0^\pm(\lambda) e_0^\pm(x, \lambda) + \overline{e_0^\pm(x, \lambda)} = \frac{2}{2 + i\beta} e_0^\mp(x, \lambda),$$

we get

$$r^\pm(\lambda) e^\pm(x, \lambda) + \overline{e^\pm(x, \lambda)} = \frac{1}{a(\lambda)} h^\mp(x, \lambda), \tag{32}_\pm$$

where

$$h^\pm(\lambda) = a(\lambda) \left[\frac{2}{2 + i\beta} e_0^\mp(x, \lambda) + l.i.m. \left(\mp \int_{\mp N}^x \Phi^\mp(x, y) e^{\pm\lambda y} dy \right) \right]. \tag{33}_\pm$$

Thus, it suffices to prove that $h^\pm(x, \lambda) = e^\pm(x, \lambda)$. If we use representations (32)_± and (33)_± for the functions $h^\pm(x, \lambda)$ and condition 2) of the theorem, then the

proof of these equalities completely coincide with the proof similar to the statement of the case $\beta = 0$ (see. [3], pp. 277-278). Therefore we don't cite it.

The theorem is proved.

Remark. Condition 4) of the theorem is essential. The function

$$r^+(\lambda) = \frac{i\beta + \frac{\gamma}{i\lambda}}{2 + i\beta + \frac{\gamma}{i\lambda}} e^{-2i\lambda a}$$

for $\gamma \neq 0$ satisfies all the conditions of the theorem, except for 4). In this case

$$K^\pm(x, t) = \begin{cases} 0, & \pm x > \pm a, \pm t > \pm x \vee \pm x < \pm a, \pm t > \pm(2a - x), \\ \frac{\gamma}{2}, & \pm x < \pm a, \pm x < \pm t < \pm(2a - x), \end{cases}$$

consequently the Jost solutions satisfy equation (3) with $q(x) \equiv 0$ and condition (4), but condition (5) is not fulfilled. If $\gamma = 0$, then condition 4) is fulfilled, and in this case, the solution of the inverse problem exists: $r^+(\lambda) = r_0^+(\lambda)$ is the right reflection factor of problem (3)-(5) with the potential $q(x) = 0$.

This work of the second author was supported by the Science Development Foundation under the President of the Republic of Azerbaijan-Grant No EIF-2011-1(3)-82/24/1

References

- [1]. Jaulent M. *On the inverse problem for the Schrodinger equation with an energy-dependent potential.* // Comptes rendus Acad. Sci. Paris, t. 280, Serie A, 1975, pp. 1467-1470.
- [2]. Jaulent M., Jean C. *The inverse problem for the on dimensional Schrodinger equation with an energy-dependent potential.* // Ann. Inst. Henri Poincare, 1976, vol. 25, No 2, pp. 105-137.
- [3]. Maksudov F.G., Huseynov G. Sh. *To the solution of the inverse scattering problem for a quadratic pencil of one-dimensional Schrodinger operators on the axis.* // DAN SSSR, 1986, vol. 289, No 1, pp. 42-46 (Russian).
- [4]. Faddeyev L.D. *Property of S-matrix of Schrodinger's one-dimensional equation.* // Trudy MIAN 1964, vol. 73, pp. 314-336 (Russian).
- [5]. Marchenko V.A. *Strum-Liouville operators and their applications.* Kiev, Naukova Dumka, 1977 (Rusiian).
- [6]. Huseynov H.M. *On continuity of reflection factor of Schrodinger one-dimensional equation.* // Diff. Uravn. 1985, vol. 21, No 11, pp. 1993-1995 (Russian).
- [7]. Savchuk A.M., Shkalikov A.A. *Sturm-Liouville operators with singular potential* // Matem. zametki. 1999, vol. 66, No 6, pp. 897-912(Russian).
- [8]. Jamshidipour A.H., Huseynov H.M. *On Jost solutions of Sturm-Liouville equations with spectral parameter in discontinuity conditions.* // Transactions of NAS of Azerb. 2010, vol. XXX, No 4, pp. 61-68.
- [9]. Jamshidipour A.H., Huseynov H.M. *Scaterring data of Sturm-Liouville operator with spectral parameter in discontinuity conditions.* // Transactions of NAS of Azerb. 2011, vol. XXXI, No 1, pp. 71-78.

[10]. Jamshidipour A.H., Huseynov H.M. *Basic equations of inverse scattering problem for Sturm-Liouville operator with spectral parameter in discontinuity conditions.* // Proceedings of IMM of NAS of Azerb. 2011, vol. XXXV, (XLIII), pp. 59-64.

Ahmad H. Jamshidipour, Hidayat M. Huseynov

Institute of Mathematics and Mechanics of NAS of Azerbaijan.

9, B.Vahabzade str., AZ 1141, Baku, Azerbaijan.

Tel.: (99412) 539 47 20 (off.).

Receivedn September 06, 2011; Revised December 15, 2011.