

MATHEMATICS

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ASYMPTOTIC BEHAVIOR OF EIGEN-VALUES OF
A BOUNDARY VALUE PROBLEM WITH
SPECTRAL PARAMETER IN THE BOUNDARY
CONDITIONS FOR THE SECOND ORDER
ELLIPTIC DIFFERENTIAL-OPERATOR EQUATION

Abstract

In this present paper we obtain the asymptotic formula for eigen values of the following boundary value problems

$$-u''(x) + Au(x) = \lambda u(x), \quad x \in (0, b),$$

$$u'(0) - \lambda u(0) = 0, \quad u'(b) + \lambda u(b) = 0,$$

where $A = A^* \geq \omega^2 I$ in H , A^{-1} is completely continuous in H , $\lambda > 0$ is a spectral parameter, H is a separable Hilbert space.

Let H be a separable Hilbert space. By $\mathcal{L}_2((0, b); H)$ ($0 < b < \infty$) we denote a set of all vector-functions $x \rightarrow u(x) : (0, b) \rightarrow H$, strongly measurable and such that $\int_0^b \|u(x)\|_H^2 dx < +\infty$. As is known $\mathcal{L}_2((0, b); H)$ is a Hilbert space with respect to the scalar product

$$(u, v)_{\mathcal{L}_2((0, b); H)} = \int_0^b (u(x), v(x))_H dx.$$

Let A be a self-adjoint positive-definite operator in H ($A = A^* \geq \omega^2 I$, $\omega > 0$, I be a unit operator in H) with domain of definition $D(A)$. Since A^{-1} is bounded in H , then

$$H(A) = \left\{ u : u \in D(A), \|u\|_{H(A)} = \|Au\|_H \right\}$$

is a Hilbert space whose norm is equivalent to the norm of the graph of the operator A .

$$W_2^2((0, b); H(A), H) = \{u : Au, u'' \in \mathcal{L}_2((0, b); H)\};$$

$\|u\|_{W_2^2((0, b); H(A), H)}^2 = \|Au\|_{\mathcal{L}_2((0, b); H)}^2 + \|u''\|_{\mathcal{L}_2((0, b); H)}^2$ is a Hilbert space [1, p.23]. In the space $\mathcal{L}_2((0, b); H)$ let's consider the boundary value problem

$$-u''(x) + Au(x) = \lambda u(x), \quad x \in (0, b), \quad (1)$$

$$u'(0) - \lambda u(0) = 0,$$

$$u'(b) + \lambda u(b) = 0, \quad (2)$$

where $A = A^* \geq \omega^2 I$ in H , A^{-1} is completely continuous in H , $\lambda > 0$ is a spectral parameter.

The goal of the paper is to study the asymptotic behavior of eigen-values of problem (1)-(2) knowing the asymptotic distribution of eigen-values of the operator A .

Note that the asymptotics of eigen-values of boundary value problems for Sturm-Liouville differential operator equation on a finite segment with the same spectral parameter in the equation and in one of the boundary conditions was studied in papers of V.I. Gorbachuk and M.A.Rybak [2], M.A.Rybak [3]. More precisely, in particular in the papers [2]-[3] it was studied the asymptotic distribution of eigen-values of the following boundary-value problem

$$-u''(x) + Au(x) = \lambda u(x), \quad x \in (0, b), \quad (1)$$

$$u'(0) + \lambda u(0) = 0, \quad (3)$$

$$u(b) = 0$$

in $\mathcal{L}_2((0, b); H) \oplus H$. It is proved that if the spectrum of the operator A is discrete, then the operator generated by the boundary value problem (1), (3) have also discrete spectrum. Eigen-values of the problem (1), (3) form two infinite sequences $\lambda_k \sim \sqrt{\mu_k}$; $\lambda_{n,k} \sim \mu_k + \frac{\pi^2}{b^2} n^2$, where $\mu_k = \mu_k(A)$ are eigen-values of the operator A .

In the paper substantially using the ideas and method of the papers [2], [3] it is proved that the problem (1)-(2) has only one set of eigen-values: $\lambda_{n,k} \sim \mu_k + \frac{\pi^2}{b^2} n^2$, where μ_k are the eigen-values of the operator.

Theorem 1. *Let $A = A^* \geq \omega^2 I$ in H and A^{-1} be completely continuous in H . Then for the eigen-values of the problem (1)-(2) it is valid the following asymptotic formula*

$$\lambda_{n,k} = \mu_k + \xi_n \quad (k = 1, 2, \dots; n = N, N + 1, \dots),$$

where $\mu_k = \mu_k(A)$ are the eigen-values of the operator A , $\xi_n \sim \frac{\pi^2}{b^2} n^2$, N is a natural number.

Proof. The eigen-elements of the operator A that correspond to eigen-values $\mu_k(A)$ we denote by φ_k . It is known that $\{\varphi_k\}$ forms an orthonormed basis in H . Then allowing for spectral expansion, for coefficients $u_k = (u, \varphi_k)$ we get the following problem

$$-u_k''(x) + (\mu_k - \lambda) u_k(x) = 0, \quad (4)$$

$$u_k'(0) - \lambda u_k(0) = 0,$$

$$u'_k(b) + \lambda u_k(b) = 0. \quad (5)$$

The general solution of ordinary differential equations (4) is of the form:

$$u_k(x) = c_1 e^{-x\sqrt{\mu_k - \lambda}} + c_2 e^{-(b-x)\sqrt{\mu_k - \lambda}}, \quad (6)$$

where c_i ($i = 1, 2$) are arbitrary constants.

Having put (6) into (5) we get a system with respect to c_i ($i = 1, 2$) whose determinant is of the form

$$K(\lambda) = -\left(\sqrt{\mu_k - \lambda} + \lambda\right)^2 + \left(\sqrt{\mu_k - \lambda} - \lambda\right)^2 e^{-2b\sqrt{\mu_k - \lambda}}.$$

So, the eigen-values of the problem (1)-(2) these zero are for the equation

$$e^{2b\sqrt{\mu_k - \lambda}} \left(\sqrt{\mu_k - \lambda} + \lambda\right)^2 - \left(\sqrt{\mu_k - \lambda} - \lambda\right)^2 = 0. \quad (7)$$

Write equation (7) in the form of a system of equations

$$e^{b\sqrt{\mu_k - \lambda}} \left(\sqrt{\mu_k - \lambda} + \lambda\right) - \left(\sqrt{\mu_k - \lambda} - \lambda\right) = 0, \quad (8)$$

$$e^{b\sqrt{\mu_k - \lambda}} \left(\sqrt{\mu_k - \lambda} + \lambda\right) + \left(\sqrt{\mu_k - \lambda} - \lambda\right) = 0. \quad (9)$$

Thus, the eigen-values of the problem (1)-(2) consist of that real $\lambda \neq \mu_k$ that even at k satisfy the equations (8)-(9).

Look for the eigen-values of the problem (1)-(2) lesser than μ_k . Put $\sqrt{\mu_k - \lambda} = y$. Equations (8) and (9) in this case are equivalent to the equations

$$e^{by} (y^2 - y - \mu_k) + (y^2 + y - \mu_k) = 0, \quad 0 < y < \sqrt{\mu_k}, \quad (10)$$

$$e^{by} (y^2 - y - \mu_k) - (y^2 + y - \mu_k) = 0, \quad 0 < y < \sqrt{\mu_k}, \quad (11)$$

respectively.

Let's prove the absence of solutions of equations (10) and (11) on the interval $(0, \sqrt{\mu_k})$. Rewrite equation (10) in the form

$$y \operatorname{sh} \frac{by}{2} + (\mu_k - y^2) \operatorname{ch} \frac{by}{2} = 0, \quad 0 < y < \sqrt{\mu_k}. \quad (12)$$

Let's consider the function $\Phi_k(y) = y \operatorname{sh} \frac{by}{2} + (\mu_k - y^2) \operatorname{ch} \frac{by}{2}$, $0 < y < \sqrt{\mu_k}$. Obviously, at each k and for all $y \in (0, \sqrt{\mu_k})$ $\Phi_k(y) > 0$. Therefore, equation (12), consequently equation (10) has no solutions on the interval k for any $(0, \sqrt{\mu_k})$.

In a similar way it is shown that equation (11) has no solutions on the interval $(0, \sqrt{\mu_k})$ for any k . Thus, equation (7) has no zeros in the case when $\lambda < \mu_k$. Now let's study that solutions of equation (7) that are greater than $\lambda > \mu_k$. In this case equations (8) and (9) are equivalent to the equations

$$e^{ibz} (z^2 + iz + \mu_k) + (-z^2 + iz - \mu_k) = 0, \quad z \in (0, +\infty), \quad (13)$$

$$e^{ibz} (z^2 + iz + \mu_k) - (-z^2 + iz - \mu_k) = 0, \quad z \in (0, +\infty), \quad (14)$$

respectively, where $z = \sqrt{\lambda - \mu_k}$.

The left hand side of equation (13) is a quasipolynomial. Applying Langer's theory on finding the asymptotics of zeros of a quasipolynomial [4, p.427] we find the asymptotics of (13). Equation (13) is equivalent to the condition

$$e^{ibz} = 1 + O\left(\frac{1}{z}\right). \quad (15)$$

The zeros of equation (15) at sufficiently great z are close (i.e. approximately equal) to the zeros of the equation

$$e^{ibz} = 1.$$

The zeros of the last equation are of the form:

$$z_n = \frac{2\pi n}{b}, \quad n = 1, 2, \dots$$

Hence for the zeros of equation (13) and hereby for the zeros of equation (8) we get the asymptotic formulae

$$\lambda_{n,k}^{(1)} = \mu_k + \xi_n, \quad k = 1, 2, \dots; \quad n = N, N+1, \dots \quad (16)$$

where $\xi_n \sim \frac{\pi^2}{b^2} (2n)^2$, N is a natural number.

Equation (14) is equivalent to the equation

$$e^{ibz} = -1 + O\left(\frac{1}{z}\right). \quad (17)$$

The zeros of equation (17) are found in a similar way. They are of the form:

$$\lambda_{n,k}^{(2)} \sim \mu_k + \frac{\pi^2}{b^2} (2n+1)^2, \quad k = 1, 2, \dots; \quad n = N, N+1, \dots \quad (18)$$

From (16) and (18) we get the asymptotic formula for the eigen-values of problem (1)-(2):

$$\lambda_{n,k} \sim \mu_k + \frac{\pi^2}{b^2} n^2, \quad k = 1, 2, \dots; \quad n = N, N+1, \dots$$

Theorem 1 is proved.

Corollary 1. Let the conditions of theorem 1 be fulfilled.

Let the eigen-values of the operator A arranged increase order satisfy the condition

$$\mu_k(A) \sim ak^\alpha \left(\lim_{k \rightarrow \infty} \frac{\mu_k(A)}{k^\alpha} = a, 0 < a, \alpha < \infty \right)$$

then the eigen-values of problem (1)-(2) have the following asymptotics

$$\lambda_m \sim dm^{\frac{2\alpha}{2+\alpha}},$$

where $0 < d < \infty$ are constants.

The proof of Corollary 1 follows from the following statement that is in the paper [5] (see also [6]).

Lemma. *Let be given the two sequences $\{\mu_k\}$ and $\{\nu_n\}$ such that $\mu_k \sim ak^\alpha, \nu_n = cn^\beta, 0 < a, c, \alpha, \beta < \infty; k = 1, 2, \dots; n = 1, 2, \dots$. Compose the sum $\mu_k + \nu_n$ with all possible k and n . We numerate the obtained numbers by increase and denote by λ_m . Then the sequences $\{\lambda_m\}$ have the asymptotics $\lambda_m \sim dm^\delta$, where*

$$\delta = \frac{\alpha\beta}{\alpha + \beta}; \quad d = \left(\frac{\alpha}{2\gamma}\right)^\delta a^{\frac{\beta}{\alpha+\beta}} c^{\frac{\alpha}{\alpha+\beta}}, \quad \gamma = \int_0^{\frac{\pi}{2}} \sin^{\frac{2}{\alpha}-1} t \cos^{1+\frac{2}{\beta}} t dt.$$

Example. In the rectangle $[0, b] \times [0, 1]$, ($0 < b < \infty$) consider the eigen-value problem

$$-\frac{\partial^2 v(x, y)}{\partial x^2} + \frac{\partial^4 v(x, y)}{\partial y^4} + \omega v(x, y) = \lambda v(x, y), \quad (19)$$

$$\frac{\partial v(0, y)}{\partial x} - \lambda v(0, y) = 0, \quad (20)$$

$$\frac{\partial v(b, y)}{\partial x} + \lambda v(b, y) = 0, \quad (21)$$

$$v_y^{(j)}(x, 0) = v_y^{(j)}(x, 1), \quad j = \overline{0, 3},$$

where $\omega > 0$ is a number.

The problem (19)-(20) is led to the form (1)-(2) where $u(x, \cdot)$ is a vector-function with values in a Hilbert space $H = \mathcal{L}_2(0, 1)$. The operator A is determined as follows:

$$D(A) = W_2^4 \left((0, 1); u^{(j)}(0) = u^{(j)}(1), j = \overline{0, 3} \right), \quad Au = \frac{d^4 u}{dy^4} + \omega u, \quad (22)$$

Obviously, the operator A determined by formula (27) is self-adjoint and for sufficiently great $\omega > 0$ is positive-definite, a A^{-1} is completely continuous in $\mathcal{L}_2(0, 1)$.

Simple calculations show that the eigen-values of the operator A are of the form: $\mu_k(A) = 16k^4 + \omega$. Then on the basis of corollary 1 the eigen-values of the problem (19)-(21) have the asymptotics: $\lambda_m \sim \text{const } m^{\frac{4}{3}}$

The author is thankful to professors S.Ya. Yakubov and B.A. Iskenderov for their useful discussions and advices.

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Received June 01, 2005; Revised September 06, 2005.

Translated by Nazirova S.H.