Ibrahim M. NABIEV, Jalala A. OSMANOVA

SOLUTION OF THE INVERSE PROBLEM FOR A SYSTEM OF EQUATIONS ON A SEGMENT

Abstract

The inverse problem is completely solved for the Dirac operator with regular boundary conditions.

1. Introduction. We consider the boundary value problems generated on the segment $[0, \pi]$ by a canonic system of Dirac equations

$$\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} y_1'(x) \\ y_2'(x) \end{pmatrix} + \begin{pmatrix} p(x) & q(x) \\ q(x) & -p(x) \end{pmatrix} \begin{pmatrix} y_1(x) \\ y_2(x) \end{pmatrix} = \lambda \begin{pmatrix} y_1(x) \\ y_2(x) \end{pmatrix}$$
(1)

and separated boundary conditions

$$y_1(0) = y_1(\pi) = 0,$$
 (2)

and also non-separated boundary conditions of the form

$$y_1(0) + \omega y_1(\pi) = \omega y_2(0) + \alpha y_1(\pi) + y_2(\pi) = 0,$$
 (3)

where p(x) and q(x) are real functions from $L_2[0, \pi]$ and ω and α are real numbers where $|\omega| \neq 1$.

We shall denote problems (1), (2) and (1), (3) by L_0 and L_1 , respectively.

At present the problems on the reconstruction of the system of equations (1) on the spectra of two boundary value problems with separated boundary conditions and also on spectral data of two problems with non-separated boundary conditions (see [1-7]) have been completely studied.

In this paper we investigate the problem on the reconstruction of problems L_0 and L_1 . The necessary and sufficient conditions were found in order that two sequences of real numbers be the spectra of such problems. The uniqueness theorem of the solution of the inverse problem is proved.

In the sequel, for the definiteness we'll assume that $|\omega| > 1$ (the case $|\omega| < 1$ is considered quite similarly). Everywhere we'll assume that k takes entire values.

Let $\begin{pmatrix} c_1(\lambda, x) \\ c_2(\lambda, x) \end{pmatrix}$ and $\begin{pmatrix} s_1(\lambda, x) \\ s_2(\lambda, x) \end{pmatrix}$ be the solutions of the system of equation (1) at initial data $c_1(\lambda, 0) - 1 = c_2(\lambda, 0) = s_1(\lambda, 0) = s_2(\lambda, 0) - 1 = 0$. Then the characteristic functions of boundary value problems L_0 and L_1 will have the form $\delta(\lambda) = s_1(\lambda, \pi)$ and

$$\Delta(\lambda) = U_{+}(\lambda) + \alpha\delta(\lambda) + 2\omega. \tag{4}$$

respectively, where

$$U_{+}(\lambda) = \omega^{2} c_{1}(\lambda, \pi) + s_{2}(\lambda, \pi).$$

$$(5)$$

It is known that [2], for the eigen-values λ_k of the problem L_0 the asymptotic formula

$$\lambda_k = k + \alpha_k, \quad \{\alpha_k\} \in l_2 \tag{6}$$

 $\frac{120}{\textit{[I.M.Nabiev, J.A.Osmanova]}}$

As in the paper [6] we can prove that the eigen-values μ_k^{\pm} of the problem L_1 are subjected to the asymptotics

$$\mu_k^{\pm} = 2k + a^{\pm} + \beta^{\pm}, \quad \{\beta_k^{\pm}\} \in l_2,$$
 (7)

where
$$a^{\pm} = \frac{2}{\pi} arctg \frac{-\alpha \pm \sqrt{\alpha^2 + (\omega^2 - 1)^2}}{(\omega - 1)^2}$$
.

2. Uniqueness theorem.

Theorem 1. The boundary value problems L_0 and L_1 are uniquely reconstructed if their spectra $\{\lambda_k\}$, $\{\mu_k^{\pm}\}$ and the sequence of signs $\sigma_k = sign[|\omega| - |s_2(\lambda_k, \pi)|]$ are known.

Proof. By (7) we have

$$\left(\frac{\omega+1}{\omega-1}\right)^2 = -\lim_{k\to\infty} tg\frac{\pi}{2}\mu_k^- \cdot tg\frac{\pi}{2}\mu_k^+,$$

$$\frac{2\alpha}{\left(\omega-1\right)^2}=-\lim_{k\to\infty}\left(tg\frac{\pi}{2}\mu_k^-+tg\frac{\pi}{2}\mu_k^+\right).$$

The coefficients ω and α of boundary conditions (3) are determined from these relations. By sequences $\{\lambda_k\}$ and $\{\mu_k^{\pm}\}$ the functions $\delta(\lambda)$ and $\Delta(\lambda)$ are reconstructed in the form of infinite products. Knowing $\alpha, \omega, \delta(\lambda)$ and $\Delta(\lambda)$ we can define $U_{+}(\lambda)$ from (4).

Denote

$$U_{-}(\lambda) = \omega^{2} c_{1}(\lambda, \pi) - s_{2}(\lambda, \pi) . \tag{8}$$

Using the identity

$$c_1(\lambda, \pi) s_2(\lambda, \pi) - c_2(\lambda, \pi) s_1(\lambda, \pi) = 1, \tag{9}$$

we can easily establish

$$U_{-}(\lambda_k) = signU_{-}(\lambda_k) \sqrt{U_{+}^2(\lambda_k) - 4\omega^2}.$$
 (10)

By (8) and (9) we have

$$U_{-}(\lambda_{k}) = \frac{\omega^{2}}{s_{2}(\lambda_{k}, \pi)} - s_{2}(\lambda_{k}, \pi) = \frac{\omega^{2} - s_{2}^{2}(\lambda_{k}, \pi)}{s_{2}(\lambda_{k}, \pi)}.$$
 (11)

It is known [2], that sign $s_2(\lambda_k, \pi) = (-1)^k$. Therefore (10) and (11) imply the equality

$$U_{-}(\lambda_{k}) = (-1)^{k} \sigma_{k} \sqrt{U_{+}^{2}(\lambda_{k}) - 4\omega^{2}}.$$

Let's consider the function

$$\varphi(\lambda) = \frac{1 - \omega^2}{4\omega^2} U_+(\lambda) + \frac{1 + \omega^2}{4\omega^2} U_-(\lambda) . \tag{12}$$

Using the representation for the functions $c_1(\lambda, \pi)$, $s_2(\lambda, \pi)$ and asymptotic formula (6) we can easily get $\{\varphi(\lambda_k)\}\in l_2$. Then by the known fact (see theorem 28 in [8]) it holds

$$\varphi(\lambda) = \delta(\lambda) \sum_{k=-\infty}^{\infty} \frac{\varphi(\lambda_k)}{(\lambda - \lambda_k) \, \delta'(\lambda_k)},$$

where
$$\varphi(\lambda_k) = \frac{1-\omega^2}{4\omega^2}U_+(\lambda_k) + \frac{1+\omega^2}{4\omega^2}(-1)^k\sigma_k\sqrt{U_+^2(\lambda_k)-4\omega^2}$$
.
Consequently, knowing $U_+(\lambda)$, $\delta(\lambda)$ and $\{\sigma_k\}$ we can uniquely define the fucn-

tion $U_{-}(\lambda)$ from (12).

It follows from (5) and (8) that $s_2(\lambda, \pi) = \frac{1}{2} [U_+(\lambda) - U_-(\lambda)]$. It is known that by the zeros of this function and by $\{\lambda_k\}$ the functions p(x) and q(x) are uniquely reconstructed.

Thus, by sequence $\{\lambda_k\}$, $\{\mu_k^{\pm}\}$ and $\{\sigma_k\}$ the coefficients of system (1) and boundary conditions (3) are uniquely reconstructed. The theorem is proved.

3. Necessary and sufficient solvability conditions of the inverse problem.

Theorem 2. In order the sequences of real numbers $\{\lambda_k\}$ and $\{\mu_k^{\pm}\}$ be the spectra of boundary value problems of the form L_0 and L_1 it is necessary and sufficient that the following conditions be fulfilled:

1) It holds asymptotic formulae (6) and (7) in which a^{\pm} are real numbers;

$$2)\dots \le \mu_k^- \le \lambda_{2k} < \mu_k^+ \le \lambda_{2k+1} \le \mu_{k+1}^- \le \dots$$
 (13)

where $\lambda_m < \lambda_{m+1}, \ m = 0, \pm 1, \pm 2...;$

$$3) \left| \Delta \left(\lambda_k \right) - 2\omega \right| \ge 2 \left| \omega \right|, \tag{14}$$

where $|\omega| > 1$, $\left(\frac{\omega+1}{\omega-1}\right)^2 = -tg\frac{\pi}{2}a^- \cdot tg\frac{\pi}{2}a^+$,

$$\Delta(z) = (\omega + 1)^2 \prod_{k=-\infty}^{\infty} \frac{(\mu_k^- - z)(\mu_k^+ - z)}{(2k + a^-)(2k + a^+)}.$$
 (15)

Proof. Necessity. Asymptotic formula (6) is in [1-2]. Formula (7) and inequality (13) are established as in theorems 2 and 3 of [6]. From equations (4) and (5) by (9) we get

$$\left|\Delta\left(\lambda_{k}\right)-2\omega\right|=\left|U_{+}\left(\lambda\right)\right|=\left|\omega^{2}c_{1}\left(\lambda,\pi\right)+\frac{1}{c_{1}\left(\lambda_{k},\pi\right)}\right|\geq2\left|\omega\right|$$

i.e. it holds inequality (14). The representation of the entire function $\Delta(z)$ in the form of infinite product (15) is established in standard way (see for example [4]).

Sufficiency. Let the conditions 1)- 3) be fulfilled.

For the function (15) it holds the representation

$$\Delta(z) = (\omega^2 + 1)\cos \pi z - \alpha \sin \pi z + f(z) + 2\omega, \tag{16}$$

[I.M.Nabiev, J.A.Osmanova]

where
$$\alpha = -\frac{(\omega - 1)^2}{2} \left(tg \frac{\pi}{2} a^- + tg \frac{\pi}{2} a^+ \right), \ f(z) = \int_{-\pi}^{\pi} \tilde{f}(t) e^{izt} dt, \ \tilde{f}(t) \in L_2[-\pi, \pi]$$

(see lemma in [6]).

By the sequence $\{\lambda_k\}$ we construct the function

$$\delta(z) = \pi (z - \lambda_0) \prod_{\substack{k = -\infty \\ k \neq 0}}^{\infty} \frac{\lambda_k - z}{k} .$$

By lemma 4 of the paper [2] this function supposes the representation of the form

$$\delta(z) = -\sin \pi z + f_1(z), \qquad (17)$$

where
$$f(z) = \int_{-\pi}^{\pi} \tilde{f}_1(t) e^{izt} dt$$
, $\tilde{f}_1(t) \in L_2[-\pi, \pi]$.

Assume

$$U_1(z) = \Delta(z) - \alpha\delta(z) - 2\omega. \tag{18}$$

By representations (16) and (17) we have

$$U_1(z) = (\omega^2 + 1)\cos \pi z + f_2(z),$$
 (19)

where $f_2(z) = f(z) - \alpha f_1(z)$. From (18) allowing for (14) we get $|U_1(\lambda_k)| \ge 2 |\omega|$, i.e. $U_1(\lambda_k) \le -2|\omega|$ or $U_1(\lambda_k) \ge 2 |\omega|$. The system of inequalities (13) shows that the signs of the terms of the sequence $\{U_1(\lambda_k)\}$ alternate. By relations (16) and (19) we get $sign U_1(\lambda_k) = (-1)^k$ for all sufficiently large values of |k|. Then there exists such a number h_k that

$$U_1(\lambda_k) = 2|\omega| (-1)^k ch h_k.$$
(20)

Consider the function

$$U_{2}(z) = \frac{\omega^{2} - 1}{\omega^{2} + 1} U_{1}(z) + \frac{4\omega^{2}}{\omega^{2} + 1} \tilde{\varphi}(z), \qquad (21)$$

where $\tilde{\varphi}(z) = \delta(z) \sum_{k=-\infty}^{\infty} \frac{\tilde{\varphi}(\lambda_k)}{(z - \lambda_k) \delta'(\lambda_k)}$,

$$\tilde{\varphi}(\lambda_{k}) = \frac{1 - \omega^{2}}{4\omega^{2}} U_{1}(\lambda_{k}) + \frac{1 + \omega^{2}}{4\omega^{2}} (-1)^{k} \sigma'_{k} \sqrt{U_{1}^{2}(\lambda_{k}) - 4\omega^{2}} =$$

$$= \frac{1 - \omega^{2}}{2|\omega|} (-1)^{k} ch \ h_{k} + \frac{1 + \omega^{2}}{2|\omega|} (-1)^{k} \sigma'_{k} |sh \ h_{k}|, \qquad (22)$$

and $\{\sigma'_k\}$ is the sequence of signs satisfying the following condition: $\sigma'_k = 0$ if $|U_1(\lambda_k)| = 2 |\omega|$ and $\sigma'_k = \pm 1$ otherwise, moreover there exists such a natural number N that $\sigma'_k = 1$ at |k| > N. By (6) and (19) $\{\tilde{\varphi}(\lambda_k)\} \in l_2$ and consequently by theorem 28 of [8] and Paley-Wiener theorem [8, p.47] it holds the representation

$$\tilde{\varphi}(z) = \int_{-\pi}^{\pi} h(t) e^{izt} dt, \quad h(t) \in L_2[-\pi, \pi] . \tag{23}$$

Since the function $s(z) = \frac{1}{2} [U_1(z) - U_2(z)]$ according to (19), (21) and (23) is representable in the form

$$s\left(z\right) = \cos \pi z + g\left(z\right) \; \left(g\left(z\right) = \int_{-\pi}^{\pi} \tilde{g}\left(t\right) e^{izt} dt, \; \; \tilde{g}\left(t\right) \in L_{2}\left[-\pi, \pi\right]\right)$$

then by lemma 4 of [2] its zeros θ_k are subjected to the asymptotics

$$\theta_k = k - \frac{1}{2} + \tau_k, \quad \{\tau_k\} \in l_2 \ .$$
 (24)

Then, allowing for (20)-(23) we have

$$s\left(\lambda_{k}\right) = \frac{1}{2}\left[U_{1}\left(\lambda_{k}\right) - U_{2}\left(\lambda_{k}\right)\right] = \left|\omega\right|\left(-1\right)^{k} ch \ h_{k}\left(1 - \sigma_{k}'\left|th \ h_{k}\right|\right)$$

and since $|th| h_k| < 1$, then $sign s(\lambda_k) = (-1)^k$. Hence, it follows that only one and by asymptotic equality (24) only one zero of the function s(z) lies on each interval $(\lambda_k, \lambda_{k+1}).$

Consequently, the zeros of the functions $\delta(z)$ and s(z) alternate

$$... < \lambda_k < \theta_{k+1} < \lambda_{k+1} < \theta_{k+2} < ...$$

Thus, the sequences $\{\lambda_k\}$ and $\{\theta_k\}$ satisfy all the conditions of theorem 2 of [2]. According to this theorem there exist the functions p(x) and q(x) such that $\delta(x)$ and s(z) are the characteristic functions of boundary value problems, generated by the system of Dirac equations with these coefficients and boundary conditions (2) and $y_1(0) = y_2(\pi) = 0$.

It is easily established that the spectrum of the constructed problem L_1 coincides with $\{\mu_k^{\pm}\}.$

The theorem is proved.

References

- [1]. Gasymov M.G., Jabiyev T.T. The solution of the inverse problem on two spectra for Dirac equation, on a finite segment. DAN Azerb. SSR, 1966, v.22, No7, pp.3-6. (Russian)
- [2]. Misyura T.V. The characteristic of spectra of periodic and antiperiodic boundary value problems generated by Dirac operator (II). Teoria funk., funct. anal. i ikh prilozh. Kharkov, 1979, issue 31, pp.102-109. (Russian)
- [3]. Nabiev I.M. The inverse problem for Dirac operator with non-separated boundary conditions. Dep. in VINITI 30.03.1987, No2281-B87.
- [4]. Levitan B.M., Sargsyan I.S. Sturm-Liouville and Dirac operators. M., Nauka, 1988. (Russian)
- [5]. Nabiev I.M. The inverse problem for Dirac operator with non-separated selfadjoint boundary conditions. Vestnik BGU, ser. fiz.-mat., 1992, No1, pp.134-138. (Russian)

$\frac{124}{\textit{[I.M.Nabiev, J.A.Osmanova]}}$

- [6]. Nabiev I.M. Solution of a class of inverse problems for the Dirac operator. Transactions of NASA, 2001, v.21, No1, pp.146-157.
- [7]. Nabiev I.M. On reconstruction of Dirac operator on the segment. Proceedings of IMM of NASA, 2003, v.XVIII, pp.97-102.
 - [8]. Levin B.Ya. Entire functions. M., MGU, 1971. (Russian)

Ibrahim M. Nabiev, Jalala A. Osmanova

Institute of Mathematics and Mechanics of NAS of Azerbaijan.

9, F.Agayev str., AZ1141, Baku, Azerbaijan.

Tel.: (99412) 439 47 20 (off.)

Received June 08, 2004; Revised October 12, 2004. Translated by Nazirova S.H.